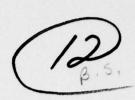
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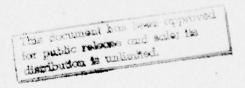


# CAVITATION INCEPTION and NUCLEI DISTRIBUTIONS JOINT ARL / CIT EXPERIMENTS



prepared by E.M. Gates, M.L. Billet, J. Katz, K.K. Ooi, J.W. Holl, A.J. Acosta

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JOINT ARL/CIT EXPERIMENTS

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California Institute of Technology Division of Engineering and Applied Science

## CAVITATION INCEPTION AND NUCLEI DISTRIBUTIONS - JOINT ARL/CIT EXPERIMENTS

#### 1. Introduction

## A. Background

Hydraulic machines (such as pumps, propellers and turbines), valves hydrofoils and any hydraulic device in which the dynamically reduced pressure falls below the vapor pressure are susceptible to cavitation. The presence of cavitation causes a loss in performance, erosion damage, and noise; hence, it is desirable to be able to predict for a given device the set of operating conditions which form the boundary between cavitating and non-cavitating states. The difficulties associated with testing or observing in operation a full scale prototype has led to the extensive use of model testing to determine cavitation behavior. To then extrapolate the model test results to full scale requires knowledge of scaling parameters and laws which adequately describe the behavior of cavitating flows.

The main parameter currently used to characterize a cavitating flow is the cavitation number  $(\sigma)$  given by:

$$\sigma = \frac{P_{\infty} - P_{V}}{\frac{1}{2} \rho V_{\infty}^{2}}$$

where  $P_{\infty}$ ,  $P_{V}$ , $\rho$ ,  $V_{\infty}$  are the liquid pressure, vapor pressure, density and velocity at infinity and the bulk temperature. The value of  $\sigma$  at which cavitation first appears is called the inception value  $(\sigma_{i})$  and the value of the cavitation number at which a cavitating condition is suppressed is called the desinent value  $(\sigma_{d})$ . A hysteresis phenomenon exists and in general  $\sigma_{d} \geq \sigma_{i}$ . However, the cavitation number is not the only parameter required to determine the inception index and despite considerable

experimental and theoretical efforts a complete set of parameters has still not been determined.

A major reason for the lack of success in determining these parameters is that there has been a great deal of confusion over the lack of reproducibility of test results from one facility to another. In an attempt to clarify this situation, the ITTC sponsored a comparative test series in which many different facilities were asked to carry out cavitation tests on a standard headform (Lindgren and Johnsson, 1966; Johnsson, 1969). The wide variation in not only the inception index but also the physical appearance of the cavitation (as is shown in Figs. 1 and 2) strongly suggested that fundamental differences in the flow conditions existed in the various facilities. As a result of the ITTC findings, there was a resurgence of interest in the area of cavitation inception and several important developments have since occurred, namely: the development of accurate methods to count and measure freestream gas bubbles and particulates (e.g. Keller, 1972; Peterson, 1972) and the discovery that viscous effects like laminar separation and boundary layer transition can play a major role in the inception process on streamlined bodies (Arakeri and Acosta, 1976). There is ample evidence that both the distribution of freestream nuclei and the type of boundary layer transition can strongly influence inception (Gates and Acosta, 1978), but the monitoring of either or both of these factors during cavitation tests is not common and hence there is little information on how these factors vary from facility to facility.

# B. Specific Objectives of Present Experiments

In the present work a comparative test series much like that of the ITTC was undertaken except that a special emphasis was placed on monitoring the flow conditions at the time of inception. Hence although the scale of the tests with regard to the number of facilities involved (3) was modest, the effort to monitor flow conditions was ambitious; for in addition to the inception index, nuclei populations were counted, schlieren observations of the viscous flow were obtained and an attempt to count cavitation "events" was also made. Further, in one facility (LTWT) nuclei populations were simultaneously measured by two techniques (holography and a counter like that of Keller's). By monitoring the different flow conditions in each facility it was hoped to

- 1. determine more clearly the influence of freestream nuclei on travelling bubble cavitation and
- 2. further observe the relation between inception and boundary layer transition.

The above objectives relate mainly to determining the fundamental processes involved in the inception of cavitation. There were also several other purposes of the study which, although as important as the ones above, are related to more practical considerations. These were:

- comparison of nuclei distributions obtained simultaneously by holography and by Keller's scattering technique.
- attempt to base the "call" of inception upon the number of cavitation events per unit time by employing a light scattering technique to count the events.

The purposes here were first to develop confidence in nuclei measurements obtained by either technique and second to investigate the possibilities of developing an arbitrary "event rate" which would define travelling bubble cavitation inception which could be applied universally.

## Equipment and Procedures

#### A. Water Tunnels

The present experiments were carried out in three water tunnels:

the High Speed Water Tunnel (HSWT), the Low Turbulence Water Tunnel (LTWT) both of which are at the California Institute of Technology and the 12-inch tunnel at the Applied Research Laboratory (ARL). The main features of each of these facilities are summarized in Table 1 and a few comments regarding the treatment of the water in each facility are made below.

## 1. 12-inch Tunnel

The circuit of the ARL facility is illustrated schematically in Fig. 3 and is described by Lehman (1959). It has no resorber, but it does have an effective deaerator with which air contents as low as 2 ppm can be obtained. Normally cavitation testing in the facility is carried out at low air contents, but during the present tests the air content was held at approximately 7 ppm. Even at this "high" air content very few freestream bubbles were visible except at inception at the low velocities (i.e. less than 30 fps).

#### 2. LTWT

The LTWT is shown schematically in Fig. 4 and the tunnel circuit is described in detail in Vanoni, et al (1950) and Gates (1977). From Fig. 4 it can be seen that the LTWT has no resorber or separate deaerator. Air is removed from the tunnel water by reducing the pressure above a free surface while the water is slowly recirculated by the tunnel pump — a minimum air content of about 7 ppm is attainable by this method. Since the maximum test section velocity is only about 25 feet per second, low test section pressures (~2 psia) are required to produce cavitation. The low pressure combined with the "relatively high air content" produces a flow in the test section which tends to have many large freestream bubbles. Also, since there is no effective way to continuously remove bubbles from the

circuit, after each test the tunnel water must be circulated for several minutes to allow the bubbles to collect and be removed from the circuit by the vacuum pump.

## 3. HSWT

A schematic drawing of the HSWT is given in Fig. 5 and the circuit details can be found in Knapp, Daily and Hammitt (1970) or more recently Ward (1976). The main feature of the facility is that it has a resorber for microbubble control and hence very few freestream bubbles are observed in the working section. Like the LTWT the HSWT has no effective system for reducing the air content. Air is removed from the tunnel water by producing super-cavitation on a test model and the sting mount. At regular intervals the tunnel is stopped and the free gas bubbles are allowed to rise and are removed. In this fashion a minimum air content of about 9 ppm may be attained after one day's operation.

## B. Nuclei Counters - Descriptions

Two optical nuclei counters were used in the present experiments.

The first is the holographic type and the second is a scattering type similar to that of Keller (1972). The scattering unit was used at all three facilities whereas the holography counter was only used at the LTWT and HSWT.

## Holographic Nuclei Counter (HNC)

Nuclei distributions were deduced from holograms of a volume of the water in the working section just ahead of the test model. The experimental apparatus and method for this counter is much the same as used by Peterson (1972), Feldberg and Shlemenson (1973) and is described in detail in Gates and Bacon (1978) and in Appendix III. Essentially it is a two step image forming process. In the first step a hologram of a sample volume of the water in the test section is recorded on a special high resolution

film by a "holocamera" (see Fig. 6). In the second step, the developed hologram is reconstructed (see Fig. 7) producing a three dimensional image of the original volume which can then be probed at the investigator's leisure. In the present tests a volume 1 cm X 1 cm X 2.5 cm near the centerline of the test section was chosen for sampling.

# Light Scattering Counter (LSC)

The light scattering nuclei counter is based on the relationship between the radius of a scattering center (R) and the scattered intensity. The scattered light intensity is also dependent on the angle of observation, the polarization angle of the scattered beam, the wavelength ( $\lambda$ ) of the scattering beam and the refraction indices of the various mediums involved. However, Keller (1972) has shown that if the scattered light is collected at an angle of  $90^{\circ}$  to the scattering beam there is a unique relationship between intensity and a non-dimensional parameter  $2\pi R/\lambda$ .

A schematic of the scattering counter is provided in Fig. 8.

First, the laser beam is expanded and collimated by a beam expander. A square aperature is used to isolate the central portion of the beam. The square beam is then focused by a long focal length lens. A small focal length lens is positioned after the focal point of the first lens so that a collimated scattering beam will result. The size of the scattering beam is determined by the size of the square aperature and the ratio of the effective focal lengths of the two lenses. Between the two lenses, there is a filter which is used to insure a scattering beam of constant intensity and an aperature positioned at the focal point to reduce reflections.

The scattered light is collected by a lens and focused by a second lens on a photomultiplier tube. A slot is positioned in front of the

photomultiplier tube to restrict the amount of scattered light collected. The size of the slot can be adjusted so that the measuring volume is cubical. The size of the control volume is determined from two considerations, namely, (1) the probability of only one scattering center being in the measuring volume at any time and (2) the characteristic dimension of the measuring volume which should be approximately three times the maximum radius of the scattering center of interest. The scattering volume in the present experiments presented a 0.76 X 0.38 mm window to the flow and was 0.76 mm deep. The output of the photomultiplier tube due to scattered light will be a pulse. This pulse goes through a signal conditioner which filters the high frequencies and adjusts the D.C. level. The pulse height processor now accepts the conditioned pulse, assesses the size of each detected pulse, and sorts the pulses into selectable categories until a predetermined total number of pulses is achieved. A display then gives the number of pulses in each category and the total amount of counting time.

During the tests at the 12-inch ARL facility the processor was programmed to divide the 0  $\rightarrow$   $\sim$  35 micrometer diameter range into 15 size categories. The number of nuclei larger than 35 micrometers diameter were indicated in channel 16. The total number of nuclei counted per sample was 1000. During the Caltech tests the size range investigated was increased to 3  $\rightarrow$  50 micrometers diameter and the number of nuclei counted per sample reduced to 250. The calibration curves giving particle diameter versus signal amplitude are provided in Appendix II.

It was previously mentioned that the scattering volume had the dimensions 0.76 mm X 0.38 mm X 0.76 mm. This is only true for the LTWT which has flat test section windows. At both the 12-inch ARL tunnel and

the HSWT it was necessary to introduce a correction for the curvature of the test section windows. This correction was determined experimentally and for the HSWT was found to be about 30 percent i.e., the sample volume was 30 percent smaller than previously stated. Since the 12-inch ARL facility has a smaller diameter than the HSWT, the correction is greater for this facility and is estimated to be about 35 percent. These corrections have been included in the reduction of the data.

## C. Flow Visualization

Thermal boundary layers on the test models were observed by schlieren photography. The particular schlieren configuration used is shown schematically in Fig. 9 and is essentially the same as that used by Arakeri (1973). Also following Arakeri, the prerequisite density gradient was produced by heating the body with internal cartridge type electric heaters. The system components and experimental technique are described in detail in Gates (1977).

At the HSWT and the ARL 12-inch tunnels correction lenses were necessary to compensate for the window curvature. Also, the windows did not have optically flat outside surfaces so it was necessary to put a thin film of glycerine between the correction lens and the tunnel window to reduce the "orange peel" effect. No such problems were encountered at the LTWT as it has flat, good quality glass test section windows.

## D. Test Models

The test model chosen for these experiments has a contour which is generated by the potential flow solution to a distributed source disk oriented normally to a uniform flow. By adjusting the source disk distribution a series of bodies can be generated each having a different minimum

pressure coefficient. The models are called Schiebe bodies (Schiebe, 1972). This model geometry was chosen for the reason that it has been shown not to have a laminar separation (van der Meulen, 1976). The significance of this is that a streamlined non-separating body typically has travelling bubble type cavitation at inception. Besides being a common type of cavitation that occurs in practice, travelling bubble cavitation is the only type of cavitation that is amenable to the concept of event or occurrence counting.

Two stainless steel Schiebe bodies with a minimum pressure coefficient of -0.75 and "final" diameters of 1.0125 and 2.025 inches were fabricated for the present experiments. The contour and pressure distribution for this particular Schiebe body are provided graphically in Fig. 10 and the information is also presented in tabulated form in Appendix 1. No quantitative measure of surface roughness was made, but each model was highly polished (a highly polished surface typically has a 0.1 X 10<sup>-6</sup>m rms finish, Beckwith and Buck, 1961). The models were supported by a one-bladed sting in the 12-inch ARL tunnel, a two-bladed sting in the LTWT and a three-bladed sting in the HSWT with the nose being at least three and normally six body diameters upstream of the sting. Misalignment from the tunnel centerline in the LTWT and HSWT is estimated to be about 0.2<sup>o</sup> but is unknown for the 12-inch ARL facility.

# E. Arrangement of Equipment at Each Facility

The physical arrangement of equipment at each facility is shown schematically in Figs. 11,13 and 15 photographically in Figs. 12, 14 and 16. In the foregrounds of Figs. !4 and 16 the light source (an argon laser) and the special filters and lenses of the LSC may be seen. In all three facilities the conditioned laser beam was reflected 90° into the test

section through the bottom window. In Fig. 12(b) the receiving optics, photomultiplier tube and a portion of the processor of the LSC are shown. The light sources and collimating lens of the schlieren system can be seen in Fig. 12(a) and the focusing lens, knife edge and camera box are shown in Fig. 16. Only in Fig. 14 can any part of the HNC be seen and that is the mirror which reflects the ruby laser pulse 90° through the test section. The holography film holder is not shown in place.

## F. Cavitation Testing Procedures

Before any experiments were carried out, the water in each facility was deaerated to reduce the number of freestream bubbles produced in the tunnel circuit. This was of particular importance in the LTWT and ARL 12-inch tunnel neither of which have a resorber. During the present tests the total air content in the 12-inch tunnel was approximately 7 ppm, in the LTWT 7-9 ppm and in the HSWT 9-10 ppm (air content levels were measured with a van Slyke blood gas analyzer).

In a typical cavitation test the following procedure was used. The water velocity in the test section was set at a specified value and kept constant. The tunnel pressure was then gradually lowered until inception occurred. Ideally at the occurrence of inception the tunnel pressures, tunnel velocity, a nuclei count, a schlieren photograph and a cavitation event rate were to be recorded. However, this was not always possible. During the tests in the ARL-12 inch tunnel it was found that the stroboscope interfered with the receiving optics of the LSC. Consequently, it was not possible to simultaneously observe inception and obtain a nuclei population. Instead, in this facility, the following method was followed: five separate observations of inception were made at a given tunnel velocity.

An average was then taken of the five tunnel pressures recorded at inception. The tunnel pressure was then reduced to this value and a "count" obtained. This procedure was repeated twice so that two counts were obtained at each tunnel speed. In the LTWT and the HSWT this difficulty was avoided by masking the strobe light sufficiently so that it did not interfere with the LSC.

In both the HSWT and the ARL-12 inch tunnel the point of inception could be approached very gradually with good pressure control. In the LTWT, however, due to the low head on the tunnel pump during an inception test the pump cavitates. Each test then had to take less than forty seconds since by that time the abundant supply of cavitation bubbles generated at the pump would reach the test section and dramatically change the free-stream conditions. Hence the tunnel pressure was reduced very rapidly until the presence of cavitation was observed. In the 12-inch tunnel and the HSWT values of cavitation desinence were also recorded.

#### F. Determination of Cavitation Inception

A standard procedure has been to observe the test body under stroboscopic lighting and to say that inception occurs when macroscopic cavities
or bubbles become visible on the model. However, this method is observerdependent and there has been a shift towards using cavitation event counters
free of human judgment. Inception is then said to occur when the "cavitation event rate" reaches a certain arbitrary value. There are problems with
event counters though, which relate mainly to the type of cavitation that
occurs and also to its location. Further, what the threshold level for
detecting an event and what the event rate at inception should be are also
open to much questioning.

During the present work an effort was made to use an optical event counter much like that of Keller's (1972) to determine inception.

Certain practical difficulties arose and it had to be abandoned and throughout all the tests inception was determined by observing the model under stroboscopic lighting. To maintain some consistent definition of inception, though, the same observers were used to "call" inception at each water tunnel.

## III. Presentation of Data

#### A. Visual Observations

## 1. Schlieren Results

In Fig. 17 an example schlieren photograph obtained on the two inch model in each facility is presented. In each case the model is seen in silhouette and the flow is from right to left. The magnification is such that the surface length shown in the photographs is approximately 10 mm. As can be seen in this figure, the Schiebe model has no laminar separation and hence transition occurs on the surface of the body. To present these observations in a more quantitative way that point on the surface at which the first noticeable disturbance occurs in the laminar boundary layer has been called the position of transition. The arc length to diameter ratio at transition, (S/D)<sub>t</sub>, has been plotted versus body Reynolds number in Figs. 18 and 19.

## 2. Photographic Results

Also during these tests 35 mm photographs of the cavitation at inception were taken to first record the type of cavitation and second to obtain an estimate of the location on the body of inception. Several different types of cavitation were observed to occur at inception and examples of

each are presented in the photographs of Fig. 20. In Fig. 20(a) "travelling bubble" or "Knappian" cavitation is presented. Here bubbles on or close to the surface grow rapidly as they approach and pass through the minimum pressure point on the body and then collapse as they continue into the region of increasing pressure. Another type of "travelling" cavitation which was observed in the ARL 12-inch tunnel is shown in Fig. 20(b). This type of cavitation will be called "travelling patch" cavitation. In this case a patch (the term patch is used to indicate that what was observed was not just a large bubble) of cavitation would appear on the model and give the impression of moving downstream along the surface of the model and then quickly disappear.

The remaining types of cavitation that occurred on the Schiebe body were attached forms of cavitation which could appear in several fashions. In one sequence a patch of cavitation would appear remain fixed in position and then disappear perhaps to reappear at another circumferential position on the model. This type is called "transient patch" cavitation and is illustrated in Fig. 20(c). In another sequence an attached cavity would form that unlike the patch type had a smooth leading edge and covered at least half of the circumference of the model — this will be referred to as a steady partially attached cavity and an example is given in Fig. 20(d). Finally, it was also observed (particularly on the 1-inch model) that the cavitation could suddenly appear as a steady attached cavity that completely circled the model circumference and had a smooth leading edge. This type has been called just that — a steady, attached cavity and is illustrated in Fig. 20(e).

As mentioned previously, the second purpose of the 35 mm photography was to attempt to locate the position of inception. This has been done and the results are presented graphically in Figs. 21 and 22 were the estimated position of inception  $(S/D)_{\hat{i}}$  has been plotted versus the body Reynolds number  $(Re_{\hat{D}})$ . In reducing the data from the photographs it was assumed that the position of inception was the further point upstream on the body when travelling bubble cavitation occurred or, in the case of an attached cavity, the leading edge of the cavity.

## B. Inception Observations

The inception index  $(\sigma_i)$  has been plotted versus the body Reynolds number for the two-inch body in Fig. 23 and for the one-inch model in Fig. 24. In each figure the type of cavitation that occurred at inception in each facility is indicated.

# 1. ARL-12 inch Tunnel

## (a) Two-inch Body

At the lowest velocity tested (30 feet/second) the following sequence of events took place: as the tunnel pressure was gradually lowered a "popping" sound became audible but no cavitation on the model was observed. A further reduction in pressure resulted in the appearance of travelling bubble type cavitation, the bubbles appearing to grow from the region of minimum pressure. Continued lowering of the tunnel pressure produced more travelling bubble cavitation and eventually a transient "patch" type of cavitation (see Fig. 20) occurred simultaneously with the travelling bubble type. As the pressure was lowered even further, the patch type became more prevalent until quite suddenly an attached, steady cavity with a laminar leading edge (see Fig. 20) appeared and there was no longer any travelling bubble cavitation.

As the freestream velocity was increased, the same sequence of

cavitation events occurred except that the difference in pressure between the first appearance of travelling bubble type inception and the first appearance of the steady attached cavity decreased. Further increases in the freestream velocity caused the transient patch type cavitation to be more prevalent at inception and with continuing increases in velocity the patch type cavitation was in its turn replaced and at velocities of 55 feet/second and higher, inception occurred as the sudden appearance of a steady, attached cavity which would cover a portion of, or all of the circumference of the body and there were no travelling bubble cavitation events.

In general, when the attached cavitation occurred it would appear very suddenly over the entire circumference and have a transparent leading edge far upstream very similar to that associated with the attached cavitation on a body having a laminar separation. However, in several instances it was observed that an attached, patch cavitation would first appear "far" downstream with a "jagged" leading edge. It would then (without any deliberate reduction in tunnel pressure) quickly "jump" upstream forming a steady, attached cavity with a smooth leading edge.

#### (b) One-inch Model

With the exception of the tests at 30 feet/second inception occurred as a steady, attached patch which would then grow circumferentially to form a cavity which completely encircled the body. At the freestream velocity of 30 feet/second travelling bubble inception occurred, but this data was taken just after the tunnel was filled. Travelling bubble events were, in general, rare on the one-inch body even though freestream conditions were similar to those during the two-inch body tests.

#### 2. LTWT

Travelling bubble type cavitation was always observed on both

test bodies during the tests in the LTWT. Lowering the tunnel pressure below the inception pressure produced profuse bubble cavitation which would eventually be replaced by a steady attached cavity supercavitation. The inception data is shown as a shaded area since an error in zeroing the pressure transducers only allows an estimate of the inception index to be made.

#### HSWT

The cavitation inception behavior in the HSWT was much the same as that in the ARL 12-inch tunnel at the higher velocities (i.e. greater than 50 feet/second). If cavitation tests were carried out immediately after filling the tunnel, inception would be of the travelling bubble type which would then very quickly be replaced by a steady attached cavity without any deliberate effort to reduce the tunnel pressure. After the tunnel water had been circulated for several minutes (approximately the time required for one complete circuit) travelling bubble events became very rare. Inception then occurred in several ways. One sequence of events was that a steady, attached patch of cavitation would first appear. This patch would then grow circumferentially until the model was completely encircled by a steady, attached cavity with a smooth leading edge. In another sequence the attached cavity would appear very suddenly without any patch type cavitation appearing. A third type of inception that occurred first took place with the appearance of an unsteady patch cavitation located "far" downstream on the test model and have a rough leading edge (see Fig. 20). This patch (or patches) would rapidly move upstream to form a steady, attached cavity.

Cavitation on the one inch model was much the same as that on the two inch model except that the sudden appearance of a steady, attached

cavity appears to be the more prevalent type of cavitation at inception.

## C. Nuclei Distributions

Nuclei populations obtained from each counter were reduced to a number density distribution function by the following approximation:

$$N\left(\frac{R_1+R_2}{2}\right) = \frac{\text{number of nuclei per unit volume with radii between } R_1 \text{ and } R_2}{(R_2-R_1)}$$

Some of these distributions are presented in Figs. 25,26 and 27 and all the nuclei counts taken in each facility are summarized in Tables II through VI.

## V. Discussion of Results

## A. Schlieren Observations

There were two reasons for making the schlieren observations: first, to monitor the viscous flow on the test body in each facility and second, to determine the relationship between inception and transition.

By observing the viscous flow in each facility it could be determined whether any differences in cavitation behavior could be partly attributed to differences in the basic flow past the test model. In Fig. 17 some schlieren photographs of transition on the two inch test model in each facility are presented. As can be seen in these photographs, transition in each facility is qualitatively the same, i.e. the transition is attached and occurs through the amplification of boundary layer waves. A more quantitative comparison is made in Figs. 18 and 19 where the position of transition has been plotted versus the body Reynolds number. It can easily be seen that all the data fall on one curve. The excellent agreement in both the qualitative and quantitative descriptions of transition in each facility lead one to conclude that the basic viscous flow about the test models in

each case is the same. Hence, any differences in cavitation behavior must be attributed to another variable. A comparison between the present observations and theory is made in Gates (1979).

## B. Relationship between Inception and Transition

During some previous tests (Gates and Acosta, 1978) carried out with a 2-inch Schiebe body in the HSWT it was found that there was a good correlation between the pressure coefficient at the position of transition and the inception index for attached forms of cavitation. It was hoped that by using the schlieren flow visualization technique that it would be possible to simultaneously observe inception and transition and hence determine if the above speculation had any validity. Unfortunately, for some practical reasons it was not possible to carry out this simultaneous observation and instead, transition information from the schlieren photographs was combined with the inception data from the 35 mm photographs and has been plotted in Figs. 21 and 22.

As can be seen in these figures most of the data clusters somewhat downstream of the position of minimum pressure and certainly does not support any correlation between inception and transition for attached cavitation. However, this information is somewhat misleading. For the attached cavitation the position of the leading edge has been plotted. This is not where inception occurs for this type of cavitation as has been described earlier. The actual position is downstream of the leading edge, but the cavity moves forward so rapidly that the true position of inception could not be determined. Hence the observations regarding attached cavitation and transition are inconclusive.

As for the travelling bubble cavitation, the implication is very

strong that it is in no way related to the position of transition but is rather controlled by the location of the minimum pressure coefficient. However, the value of the inception index is not only determined by the minimum pressure coefficient but also by the nuclei population as we will see later.

## C. Nuclei Populations

Nuclei populations in the ARL-12 inch tunnel were measured with the L.S.C. only. A few representative populations have been reduced to number density functions and are presented in Fig. 25 and all the distributions are summarized in Table II. As can be seen from the table, the nuclei density appears to be relatively constant for the range of velocities and pressures covered. From this one might deduce that the nuclei are mainly particulates, but as shall be seen this is not believed to be the case.

Nuclei populations in the HSWT were measured using both optical counters but on separate occasions approximately 11 months apart in time. Again a sample population from each counter has been reduced to a number density function and is shown in Fig. 26. The counts have also been summarized in Table III and IV.

The following observations were made with the LSC counter: at any given velocity the nuclei population was found to be independent of pressure. If the velocity was gradually increased, it was observed that a large number of small nuclei were produced. Also, if the velocity and pressure were held constant, the number of small particles (>10 micrometers diameter) increased whereas the number of large particles decreased with time. From these observations it was deduced that the nuclei in the HSWT are predominantly particulates which are believed to be particles of rust

and chips of paint which are continuously being removed from the tunnel walls. This conclusion agrees with the holography observations in which so few bubbles were observed that no estimate of a distribution could be made, i.e. the distribution presented in Fig. 26 for the HSWT is for particulates.

In the LTWT it was possible to simultaneously measure nuclei populations using both optical counters. A comparison of some distributions deduced from these measurements is made in Fig. 27 and in Table V all the available populations are compared in tabular form. The limited range of velocity and lack of precise pressure control in the LTWT made it difficult to vary tunnel conditions significantly to deduce from the LSC the relative populations of particulates and gas bubbles. From the HNC results, however, it is possible to classify the nuclei and in Table VI the holography results of Table V have been presented again but with the particulate and bubble populations separated. No distinction between particulate and bubble populations for nuclei less than 20 micrometers diameter is possible as background noise destroys the resolution. Nuclei between 20-50 micrometers diameter are mainly particulates, between 50-100 micrometers there are approximately equal numbers of each and above 100 micrometers diameter the nuclei are believed to be essentially all bubbles.

It can readily be seen from these results that there is a substantial difference in the populations measured by each nuclei counter. This discrepancy is discussed in more detail in Billet and Gates (1979), but no conclusion that could definitely explain the difference was drawn. This is a particularly important problem to be resolved if any reliable comparison between the performance of various cavitation facilities and measurement devices are to be made.

## D. Nuclei and Inception

First consider the results for the 2-inch Schiebe body in the ARL 12-inch tunnel. As the freestream velocity was increased two trends in the inception behavior were noted. One, the inception index decreased and second, the type of cavitation at inception gradually changed from travelling bubble to attached cavitation. If the nuclei population remained constant, it would be expected that the inception index should rise since the number of encounters between model and nuclei will increase with velocity. The speculation is that at the higher freestream velocities the static tunnel pressure is high at inception and hence one might expect fewer freestream gas bubbles to be present. And, if it is assumed that particulates do not influence inception, then a decreasing inception index would be expected. The trend to an attached form of cavitation tends to support this speculation. For, it was found in the HSWT where there were very few freestream gas bubbles that attached forms of cavitation occurred whereas in the extremely bubbly flow of the LTWT only travelling bubble cavitation was observed.

Further supporting evidence comes from the inception observations upon the 1-inch Schiebe body in the ARL facility. The smaller model will "see" fewer nuclei than the 2-inch model. The main type of cavitation observed on the 1-inch body at inception was the attached type. The implication being that at a given velocity there may be enough nuclei to produce travelling bubble inception on the 2-inch model whereas there are too few on the 1-inch model and attached cavities appear at inception.

Cavitation inception in the LTWT was always of the travelling bubble type and even though there is some question as to the absolute value

of the inception index it tends to be lower than the values observed in the ARL 12-inch tunnel. This is surprising in that the LSC measured approximately an order of magnitude more nuclei/cc in the LTWT than in the ARL facility.

In the HSWT only attached forms of cavitation occurred and surprisingly for a given freestream velocity the inception index was the same for both the 1-inch and 2-inch bodies. (It may be noted here that if cavitation depends upon the number of nuclei "swept out" by the body that a smaller body should have a smaller value of  $\sigma_i$  for a constant concentration of nuclei per unit volume.) Further, as can be seen in Fig. 23, there is a suggestion of a correlation between the pressure coefficient at transition and the inception index although, this is not supported by observations of the position of transition in the present tests.

In all three facilities then, there are substantial differences in the type of cavitation at inception and the inception index. Since the viscous flow past each model has been shown to be the same in each facility it is apparent that these discrepancies are a consequence of different nuclei populations. The observations in the HSWT indicate first that particulates (at least those in the HSWT) are not sites for nucleation and second that when extremely few freestream gas bubbles are present that a simple explanation based on cavitation event encounters cannot be used. Instead it appears that some interaction between nuclei and transition is responsible for inception. At the other end of the spectrum in the LTWT where there are very many freestream bubbles, travelling bubble inception will occur and the inception index is then controlled by the bubble population.

These cavitation observations are somewhat inconclusive. Even though the LSC measures approximately an order of magnitude fewer nuclei in the ARL tunnel than the LTWT, travelling bubble inception occurs at a higher inception index than in the LTWT. Yet at the same time there appears to be a scarcity of nuclei in the ARL 12-inch water tunnel so that an attached cavity forms on the 2-inch model at the higher velocities and on the 1-inch model over the entire velocity range. (However, it may be pointed out that if the attached cavitation is a consequence of surface relative roughness we would expect the smaller body to be the rougher and hence that attached cavitation would form more readily then.)

## VI. Concluding Remarks

From the introductory remarks it should be recalled that the main objectives of this investigation were to further pursue the questions of the relationships between transition and nuclei populations with inception. To carry out this work a non-separating test model was chosen under the assumption that it would exhibit travelling bubble type cavitation at inception. Surprisingly, this type of inception was only predominant in the extremely bubbly flow of the LTWT whereas in the HSWT and the ARL-12 inch tunnel attached forms of inception occurred.

Schlieren observations of the viscous flow about the model demonstrated that it was unaltered from facility to facility and hence it is believed that differences in cavitation behavior are attributable to nuclei populations and in particular the gas bubble fraction of these populations. Hence although the LSC indicates that the populations in the ARL tunnel and the HSWT are about the same, it is speculated that the bubble fraction of the distribution is smaller in the resorber facility i.e.

the HSWT. When very few gas bubbles are present as in the HSWT the cavitation at inception tends to be of the attached type and the lack of diameter effect here implies the concept of encounter rates is not applicable to defining inception. Whereas in extremely bubbly flows like that of the LTWT the encounter rate seems to be appropriate.

The consequence of these observations then is to re-emphasize the need for accurate measurements of nuclei popultions and in particular to be able to distinguish between particulates and bubbles. The rather significant difference between counters must then be resolved before any reliable comparison of cavitation behavior from one facility or environment to another may be made.

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TABLE I WATER TUNNEL CHARACTERISTICS

	12-inch ARL	HSWT	LTWT		
Type of Circuit	Closed	Closed	Closed		
Working Section	Cylindrical 12-inch diameter	Cylindrical 14-inch dia. 48-inches long	Square 12 inch X 12 inch 96 inch long		
Maximum Velocity (fps)	80	100	25		
Pressure Range: Min(atm) Max(atm)	0 4	0.1 7.5	0.1 1.0		
Resorber	No	Yes	No		
Typical Air Content during tests mole/mole	4	9	7		
Freestream Turbulence Level (percent)	0.15	0.2	0.05 → 3.6		

Table 11
Nuclei Distributions in ARL 12 inch Tunnel
as Measured by LSC

	Cumulative											-28												
	Cumul		3.80	2.53	14.1	18.2	17.2	14.9	12.2	14.3	14.1	12.2	5.6	2.7		32.3	0.04		11.2	16.1	11.9	11.9	7.6	
<b>9</b>	> 21.5		0.17	0.13	0.23	0.24	0.21	0.10	0.11	0.14	0.14	0.10	0.03	0.02		0.55	09.0		1.17	1.42	1.44	0.88	0.84	
Nuclei Distribution (no./cc) Diameter (micrometers)	19-21.5		0.03	0.03	0.08	0.09	0.02	0.03	0.04	0.07	80.0	0.05	0.01	0.01	book.	0.03	40.0		0.38	0.47	0.30	0.48	0.43	
stributi eter (mi	16-19		90.0	0.05	0.14	0.11	0.07	0.12	0.11	0.11	0.10	90.0	0.01	0.02		0.29	91.0		0.45	95.0	0.52	0.40	0.31	
Nuclei Di Diam	12.5-16	en 2/28/78	91.0	0.10	0.25	0.27	0.38	0.37	0.33	0.27	0.32	0.28	0.03	40.0	in original data	0.87	96.0	en 3/2/78	0.21	0.47	0.33	0.45	91.0	
	10-12.5	Following Data Taken 2/28/78	0.22	0.13	0.65	97.0	0.62	0.45	0.48	94.0	0.45	0.23	40.0	0.03	as suspect	0.87	1.28	Following Data Taken 3/2/78	0.38	99.0	0.45	0.30	0.37	
	6.5-10	Followin	0.33	0.25	1.27	1.25	1.09	1.01	9.65	1.07	0.62	0.62	0.11	0.12		5.06	3.04	Followin	0.27	44.0	0.18	0.29	0.26	
3	3-6.5		0.34	0.32	1.21	1.60	0.98	0.87	0.72	0.91	98.0	0.63	0.15	0.13	two counts noted	2.65	2.96		0.80	1.05	69.0	0.70	0.55	
	1-3		2.51	1.52	10.25	13.85	13.88	12.00	9.77	11.24	11.54	10.22	2.23	2.31	Above	24.94	30.96		7.57	90'11	8.02	8.44	6.79	
	Volume	20	263	395	11	55	28	19	82	70	11	82	383	373		31	25		89	62	84	84	103	
	ь		0.53-0.65	0.73-0.79	0.68-0.70	0.64-0.69	0.63	0.60-0.63	0.56-0.57	0.56	0.54	0.51-0.53	0.50-0.51	0.50-0.51		0.50-0.51	0.51-0.53		0.73-0.76	0.78-0.80	0.57-0.60	0.59-0.61	0.59-0.60	
	Velocity	ļ.	30.3	30.3	35.5	35.5	40.2	40.2	45.3	45.3	50.4	4.05	55.3	55.3		4.09	4.09		29.1	28.6	36.0	36.1	39.9	

Table 11 continued

Cumulative		3.0	3.0		5.0	5.0	3.6
>21.5		0.21	0.24		0.31	0.40	0.32
19-21.6		0.10	0.09	lodel	0.11	0.11	0.09
16-19	80	0.13	0.11	o clean m	0.17	0.14	0.13 0.09
12.5-16	cen 3/3/7	0.08 0.13	0.10	d opened t	0.11	90.0	0.07
3-6.5 6.5-10 10.12.5 12.5-16 16-19 19-21.6 > 21.5	Following Data Taken 3/3/78	0.11	0.0 0.10 0.10 0.09 0.3	topped and	0.11	0.12	0.10
6.5-10	Followin	0.07	90.0	Tunnel s	0.10	0.12	0.07
3-6.5		0.18	0.19		0.27	0.33	0.20
<u>.</u>		2.09	2.14		3.84	3.67	2.63
Volume Sampled cc		336	330		200	201	772
ь		0.43-0.49	0.46-0.48		0.44-046	0.47-0.48	0.42
Velocity fps		45.3	45.3		50.4	55.3	60.5

TABLE III

Nuclei Distributions in HSWT
as Measured by LSC

Nuclei Distribution (no./cc) Nuclei Diameter (micrometers)

Velocity fps	σ	Volume Sampled cc	< 11	11-21.5	21.5-50	> 50	Cumulative
20.83	4.75	224	3.60	0.64	0.06		4.46
20.76	0.96	244	3.60	0.24	0.08		4.10
20.78	4.77	222	4.16	0.28	0.06	-	4.50
			Rest 10	minutes			
31.14	2.12	154	5.88	0.28	0.10	0.24	6.50
31.25	0.43	304	2.84	0.40	0.04	0.02	3.30
31.28	2.10	234	4.12	0.36	0.12	0.02	4.62
			Rest 10	minutes			
46.82	0.94	164	4.58	0.60	0.04		5.22
46.72	0.51	66	12.78	0.54	0.18	1.64	15.14
46.93	0.94	56	13.42	0.92	0.08	2.34	17.84
			Rest 10	minutes			
62.70	0.66	21	44.20	1.34		2.10	47.64
62.82	0.54	22.2	43.24	1.62	0.18		45.04
62.76	0.66	13.2	76.66	0.90		0.60	78.16
			Rest 10	minutes			
62.76	0.66	12.2	63.60	0.66	<u>.</u>	17.70	81.96
			1 minute	e			
62.76	0.66	13.4	62.98	0.60	0.30	10.74	74.62
			1 minute	e			
62.76	0.66	14.4	65.56	0.84	_	3.06	69.46
			6 minute	es			
62.76	0.66	14.4	67.50	1.12	0.28	0.56	69.46
			2 minute				
62.76	0.66	15.6	62.56	1.02		0.52	64.10
			1 minute				

Table III continued

5 minutes 49.20 0.80 50.00  5 minutes  62.93 0.41 21.16 45.94 0.94 0.18 0.18 47.24  5 minutes  62.93 0.41 16.70 57.72 0.72 0.48 0.72 59.64  Run 10 minutes at Atmos. Press.  30.60 2.20 57.40 16.24 1.04 0.14 - 17.42  30.60 0.45 49.28 18.66 1.46 0.16 - 20.30  Tunnel Rest for 30 minutes  46.68 0.94 43.78 21.66 1.10 0.10 - 22.86  46.90 0.49 31.54 29.16 1.14 - 1.40 31.70  46.91 0.93 33.20 29.16 0.72 0.24 - 30.12  Tunnel Rested Overnight  31.62 0.93 292 2.94 0.32 0.16 - 3.42  31.62 0.93 310 2.52 0.56 0.12 0.02 3.22  Cavitation Testing Starts  31.23 0.27 236 3.46 0.64 0.14 - 4.24  30.89 0.25 272 3.00 0.48 0.18 0.02 3.68  41.79 0.36 272 3.14 0.44 0.08 0.02 3.68  41.79 0.36 272 3.14 0.44 0.08 0.02 3.68  41.79 0.36 272 3.14 0.44 0.08 0.02 3.68  52.60 0.45 284 3.12 0.14 0.04 0.22 3.52  62.28 0.47 109.12 8.32 0.62 0.10 0.10 9.14  63.04 0.60 53.56 10.76 0.82 0.14 6.94 18.66  30.49 0.41 724 1.10 0.22 0.06 - 1.38  41.32 0.33 608 1.24 0.36 0.04 0.02 1.66	Velocity fps	σ	Volume Sampled cc	< 11	11-21.5	21,5-50	> 50	Cumulative
62.93	62.3	0.41	20.0			•	-	50.00
5 minutes 62.93				5 minute	es			
62.93	62.93	0.41	21.16	45.94	0.94	0.18	0.18	47.24
Run 10 minutes at Atmos. Press.  30.60				5 minute	es			
30.60	62.93	0.41	16.70	57.72	0.72	0.48	0.72	59.64
30.60 0.45 49.28 18.66 1.46 0.16 - 20.30  Tunnel Rest for 30 minutes  46.68 0.94 43.78 21.66 1.10 0.10 - 22.86  46.90 0.49 31.54 29.16 1.14 - 1.40 31.70  46.91 0.93 33.20 29.16 0.72 0.24 - 30.12  Tunnel Rested Overnight  31.62 0.93 292 2.94 0.32 0.16 - 3.42  31.62 0.93 310 2.52 0.56 0.12 0.02 3.22  Cavitation Testing Starts  31.23 0.27 236 3.46 0.64 0.14 - 4.24  30.89 0.25 272 3.00 0.48 0.18 0.02 3.68  41.79 0.36 272 3.14 0.44 0.08 0.02 3.68  41.79 0.36 272 3.14 0.44 0.08 0.02 3.68  41.79 0.36 272 3.14 0.44 0.08 0.02 3.68  41.79 0.36 272 3.14 0.44 0.08 0.02 3.68  41.79 0.36 272 3.14 0.44 0.08 0.02 3.68  41.79 0.36 272 3.14 0.44 0.08 0.02 3.68  41.79 0.36 272 3.14 0.44 0.08 0.02 3.68  41.79 0.36 272 3.14 0.44 0.08 0.02 3.68  41.79 0.36 272 3.14 0.44 0.08 0.02 3.68  41.79 0.36 272 3.14 0.44 0.08 0.02 3.68  41.79 0.36 272 3.14 0.44 0.08 0.02 3.68  41.79 0.36 272 3.14 0.44 0.08 0.02 3.68			Run 10 m	inutes at	Atmos. Pres	s.		
Tunnel Rest for 30 minutes 46.68	30.60	2.20	57.40	16.24	1.04	0.14	-	17.42
46.68	30.60	0.45	49.28	18.66	1.46	0.16		20.30
46.90				Tunnel F	Rest for 30	minutes		
46.91 0.93 33.20 29.16 0.72 0.24 - 30.12  Tunnel Rested Overnight  31.62 0.93 292 2.94 0.32 0.16 - 3.42  31.62 0.93 310 2.52 0.56 0.12 0.02 3.22  Cavitation Testing Starts  31.23 0.27 236 3.46 0.64 0.14 - 4.24  30.89 0.25 272 3.00 0.48 0.18 0.02 3.68  41.79 0.36 272 3.14 0.44 0.08 0.02 3.68  52.60 0.45 284 3.12 0.14 0.04 0.22 3.52  62.28 0.47 109.12 8.32 0.62 0.10 0.10 9.14  63.04 0.60 53.56 10.76 0.82 0.14 6.94 18.66  30.49 0.41 724 1.10 0.22 0.06 - 1.38  41.32 0.33 608 1.24 0.36 0.04 0.02 1.66	46.68	0.94	43.78	21.66	1.10	0.10	•	22.86
Tunnel Rested Overnight  31.62	46.90	0.49	31.54	29.16	1.14	-	1.40	31.70
31.62	46.91	0.93	33.20	29.16	0.72	0.24	-	30.12
31.62 0.93 310 2.52 0.56 0.12 0.02 3.22 Cavitation Testing Starts  31.23 0.27 236 3.46 0.64 0.14 - 4.24 30.89 0.25 272 3.00 0.48 0.18 0.02 3.68 41.79 0.36 272 3.14 0.44 0.08 0.02 3.68 52.60 0.45 284 3.12 0.14 0.04 0.22 3.52 62.28 0.47 109.12 8.32 0.62 0.10 0.10 9.14 63.04 0.60 53.56 10.76 0.82 0.14 6.94 18.66 30.49 0.41 724 1.10 0.22 0.06 - 1.38 41.32 0.33 608 1.24 0.36 0.04 0.02 1.66				Tunnel F	Rested Overn	ight		
Cavitation Testing Starts  31.23	31.62	0.93	292	2.94	0.32	0.16	-	3.42
31.23       0.27       236       3.46       0.64       0.14       -       4.24         30.89       0.25       272       3.00       0.48       0.18       0.02       3.68         41.79       0.36       272       3.14       0.44       0.08       0.02       3.68         52.60       0.45       284       3.12       0.14       0.04       0.22       3.52         62.28       0.47       109.12       8.32       0.62       0.10       0.10       9.14         63.04       0.60       53.56       10.76       0.82       0.14       6.94       18.66         30.49       0.41       724       1.10       0.22       0.06       -       1.38         41.32       0.33       608       1.24       0.36       0.04       0.02       1.66	31.62	0.93	310	2.52	0.56	0.12	0.02	3.22
30.89       0.25       272       3.00       0.48       0.18       0.02       3.68         41.79       0.36       272       3.14       0.44       0.08       0.02       3.68         52.60       0.45       284       3.12       0.14       0.04       0.22       3.52         62.28       0.47       109.12       8.32       0.62       0.10       0.10       9.14         63.04       0.60       53.56       10.76       0.82       0.14       6.94       18.66         30.49       0.41       724       1.10       0.22       0.06       -       1.38         41.32       0.33       608       1.24       0.36       0.04       0.02       1.66				Cavitati	ion Testing	Starts		
41.79       0.36       272       3.14       0.44       0.08       0.02       3.68         52.60       0.45       284       3.12       0.14       0.04       0.22       3.52         62.28       0.47       109.12       8.32       0.62       0.10       0.10       9.14         63.04       0.60       53.56       10.76       0.82       0.14       6.94       18.66         30.49       0.41       724       1.10       0.22       0.06       -       1.38         41.32       0.33       608       1.24       0.36       0.04       0.02       1.66	31.23	0.27	236	3.46	0.64	0.14	•	4.24
52.60       0.45       284       3.12       0.14       0.04       0.22       3.52         62.28       0.47       109.12       8.32       0.62       0.10       0.10       9.14         63.04       0.60       53.56       10.76       0.82       0.14       6.94       18.66         30.49       0.41       724       1.10       0.22       0.06       -       1.38         41.32       0.33       608       1.24       0.36       0.04       0.02       1.66	30.89	0.25	272	3.00	0.48	0.18	0.02	3.68
62.28       0.47       109.12       8.32       0.62       0.10       0.10       9.14         63.04       0.60       53.56       10.76       0.82       0.14       6.94       18.66         30.49       0.41       724       1.10       0.22       0.06       -       1.38         41.32       0.33       608       1.24       0.36       0.04       0.02       1.66	41.79	0.36	272	3.14	0.44	0.08	0.02	3.68
63.04	52.60	0.45	284	3.12	0.14	0.04	0.22	3.52
30.49	62.28	0.47	109.12	8.32	0.62	0.10	0.10	9.14
41.32 0.33 608 1.24 0.36 0.04 0.02 1.66	63.04	0.60	53.56	10.76	0.82	0.14	6.94	18.66
	30.49	0.41	724	1.10	0.22	0.06	•	1.38
E2 56 0 hE 220 2 62 0 16 0 0h 0 5h h 26	41.32	0.33	608	1.24	0.36	0.04	0.02	1.66
52.50 0.45 250 5.02 0.10 0.04 0.54 4.50	52.56	0.45	230	3.62	0.16	0.04	0.54	4.36

Table 1V Nuclei Populations in HSWT as Measured by Holography (Katz 1978)

Nuclei Distribution (no./cc) Nuclei Diameter (micrometers)

σ	< 40	40-100	100-150	150-200	> 200	Cumulative
0.238	40.4	16.8	4.0	0.4	0.8	62.40
0.257	44.4	8.8	3.6	2.8	1.2	60.80
.223	32.0	14.0	5.6	1.6	2.0	55.2
.230	26.8	7.2	1.2	0.4	0.4	36.0
0.410	35.6	15.2	6.4	1.6	0.8	58.8
.469	42.0	21.2	2.4	0.4	•	66.0
.433	41.2	14.4	5.6	1.6	0.4	63.2
.431	24.4	8.4	2.0	0.4	0.4	35.6
.464	30.0	11.2	2.8	1.6	0.8	46.4
.615	14.0	4.4	-	-	0.4	18.8
0.647	38.0	14.8	2.0	1.2	•	56.0
0.610	26.0	14.8	3.2	1.2	0.4	45.6
0.613	37.6	14.8	3.6	0.8	0.4	57.2
727	28.4	6.4	0.8	1-	•	35.6
737	27.2	8.8	2.4	0.4	0.4	39.2

TABLE V

Comparison of Nuclei Measurements Obtained in LTWT

Nuclei Distribution (number/cc) Nuclei Diameter (micrometers)

ring						-33	3-														
ative Light Scattering	•	٠			•	•	•	87.8	81.0	88.4	0.99	88.0	212.0	140.0	202.0	175.6	490.0	510.0	136.0	128.0	176.0
Cumulative Hplo. Lig Sca	315	114	394	411	372	318	422	•	•	ı	•	197	604		•	•	٠	٠	•	914	394
Light Scattering	•		•		ı	•	1	1.40	49.0	2.48	90.1	1.0	8.0		1.60	1.40	2.06	4.08	2.18	9.0	•
>50 Holo.	8.0	7.2	2.4	1.2	2.8	2.8	111		•	•	•	1.2	3.2	•				•	•	9.1	9.1
50 Light Scattering		•		•	1	•	ı	3.52	4.22	96.4	3.96	9.4	4.2	0.56	7.22	2.10	2.68	12.24	0.54	3.6	2.8
20-50 Holo.	82.8	126	126	911	84	79	110	•			•	39	911	•			•	•		901	18
20 Light Scattering	•	•			1	•	•	82.80	76.20	81.0	0.19	82.0	208.0	139.4	193.6	172.0	486.8	0.464	133.2	126.0	172.0
< 20 Holo.	224	338	597	294	285	239	308	ı	٠		•	221	290	•	ı	•	1	•	•	308	312
σ (est)	94.0	0.45	0.39	0.38	0.39	0.40	0.39	35.78	35.34	32.92	25.49	25.83	3.19	3.28	2.13	1.99	1.42	1.41	0.45	0.47	0.47
Velocity fps	23.5	23.3	23.3	23.3	23.5	23.4	23.5	4.5	4.5	4.5	4.2	4.2	20.4	20.2	20.1	20.1	20.0	20.1	22.5	22.6	22.2
Hologram Number	-	4	5	9	7	8	6	12	13	14	15	91	17	18	19	20	21	22	23	24	25

	gu .												-34	-	
ıla	Light Scattering	136.8	146.0	149.0	139.0	186.2	146.8	9.99	168.0	148.0	180.0	160.0	152.8	166.8	175.8
	Holo.						•		273	310		•	٠	٠	•
	Light Scattering		9.0	9.0	•	1.4	9.0	•	9.0	•	8.0	9.0	•	•	0.7
>50	Holo.	•		•	•	•	•		2.8	4.4	•		•		•
-50	Light Scattering	2.74	1.8	3.0	2.8	0.9	•	9.1	9.0	9.1	5.0	3.2	2.4	9.0	0.7
	Holo.	•	•	•	•		•	•	99	49	•			•	ı
•	Light Scattering	134.0	144.0	145.4	136.2	178.8	146.2	65.	0.991	146.0	174.8	156.2	150.4	166.2	174.4
< 20	Holo.								204	242	•	ı	•	•	•
ь	(est.)	0.48	0.51	0.50	0.47	09.0	0.44	0.52	0.54	0.51	0.51	0.48	0.45	0.43	0.47
Velocity	Fps	22.6	22.6	22.7	22.7	18.5	20.7	20.9	20.3	21.6	21.0	20.9	23.1	23.5	23.3
Hologram	Number	26	27	28	29	30	31	33	34	35	36	37	38	39	40

Table VI
Particulate and Bubble Distributions in the LTWT

Nuclei Distribution (number/cc) Nuclei Diameter (micrometers)

Hologram	20-	50	50-100				
Number	Part.	Bubbles	Part.	Bubbles			
1	82.0	0.8	8.0	-			
4	124.8	1.2	6.0	1.2			
5	122.4	3.2	2.4	•			
6	113.6	2.4	0.4	0.8			
7	70.0	14.4	-	2.8			
8	69.2	10.0	1.2	1.6			
9	104.0	6.0	0.4	4.0			
17	98.0	17.6	0.4	2.8			
24	92.8	12.8	0.4	1.2			
25	70.0	10.8	0.4	1.2			

# FIGURE CAPTIONS

- Figure 1: Results of a comparative cavitation inception test on a modified ellipsoidal headform sponsored by the ITTC.
- Figure 2: Photographs of incipient cavitation on the ITTC headform in various facilities.
- Figure 3: A schematic drawing of the ARL 12-inch water tunnel.
- Figure 4: A schematic drawing of the LTWT circuit.
- Figure 5: A schematic drawing of the HSWT circuit.
- Figure 6: Schematic drawing illustrating the main components of the holocamera.
  - (a) etalon
  - (b) iris
  - (c) dye cell
  - (d) ruby rod flash lamp assembly
  - (e) iris
  - (f) dielectric mirror
  - (g) beam splitter
  - (h) neutral density filter
  - (i) beam expander lens
  - (j) 25µ pinhole
  - (k) collimating lens
  - (1) front surface mirror
  - (m) pin diode
  - (n) film pack
- Figure 7: A schematic representation of the system to reconstruct the holograms.

- Figure 8: A schematic drawing of the basic components of the light scattering counter.
- Figure 9: A schematic drawing of the schlieren flow visualization system.
- Figure 10: The pressure coefficient versus arc length for the C  $_{\rm P_{min}}$  = -0.75 Schiebe body.
- Figure 11: A schematic representation of the arrangement of equipment at the ARL 12-inch tunnel.
- Figure 12: Two photographs illustrating the arrangement of equipment at the ARL 12-inch tunnel.
- Figure 13: A schematic drawing of the experimental configuration at the LTWT.
- Figure 14: Aphotograph of the actual equipment used at the LTWT.
- Figure 15: A schematic drawing illustrating the equipment arrangement at the HSWT.
- Figure 16: A photograph of the equipment arrangement at the HSWT.
- Figure 17: A sequence of schlieren photographs illustrating transition on the 2 inch Schiebe body in each facility. In each photograph the model is seen in silhouette with the flow from right to left and the arc length is approximately 10mm.
  - (a) ARL 12-inch tunnel

$$V = 35.1 \text{fps } S/D_t = 0.67$$

(b) HSWT

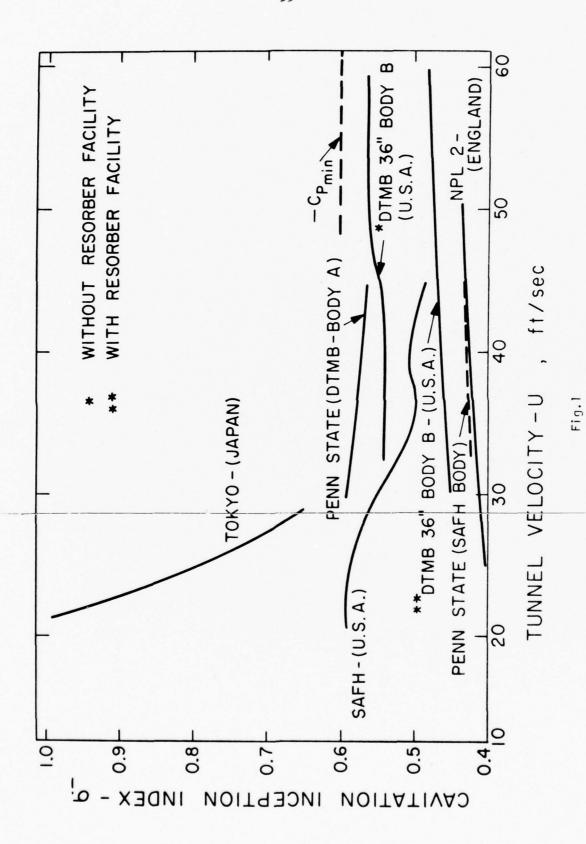
$$V = 30.0 fps S/D_{c} = 0.72$$

(c) LTWT

$$V = 20.7 fps$$

Figure 18: A summary of the averaged schlieren observations of transition on the 2-inch diameter Schiebe model.

- Figure 19: A summary of the averaged schlieren observations of transition on the one inch diameter Schiebe model.
- Figure 20: A sequence of photographs showing the several types of incipient cavitation that occur on the Schiebe models.
  - (a) travelling bubble LTWT, V = 24.0 fps,  $\sigma = 0.53$
  - (b) travelling patch ARL 12 inch tunnel and ARL model V = 60 fps,  $\sigma = 0.365$
  - (c) patch type cavitation HSWT, V = 41.5 fps,  $\sigma$  = 0.43
  - (d) partially attached and patch type cavitation HSWT V = 41.6 fps,  $\sigma = .40$
  - (e) steady, fully attached cavitation HSWT, V = 41.6 fps  $\sigma = 0.39$
- Figure 21: A plot of the estimated position of inception versus body Reynolds number for the two inch diameter Schiebe body.
- Figure 22: A graph of the estimated position of cavitation inception versus
- Figure 23: A summary of the cavitation inception data for the two inch diameter Schiebe model in each facility.
- Figure 24: A summary of the cavitation inception data for the one inch diameter Schiebe model in each facility.
- Figure 25: Several nuclei distribution functions calculated from populations obtained with the LSC in the ARL 12-inch tunnel.
- Figure 26: A comparison of nuclei populations measured in the HSWT with the LSC and the HNC at a time interval of about 11 months.
- Figure 27: A comparison of nuclei populations measured simultaneously in the LTWT by the two optical counters.



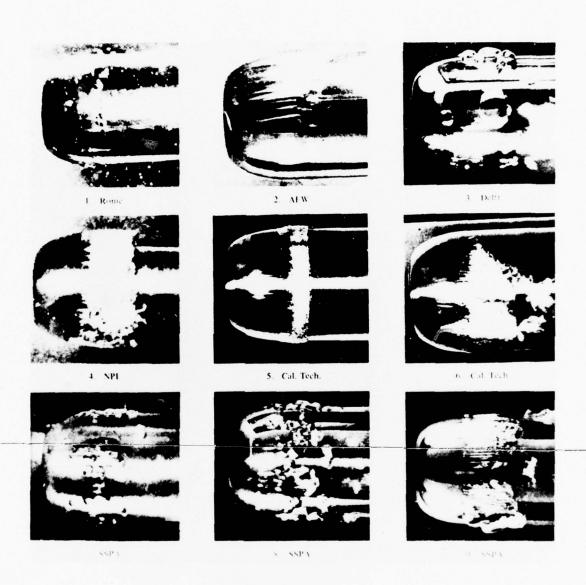


Fig. 2

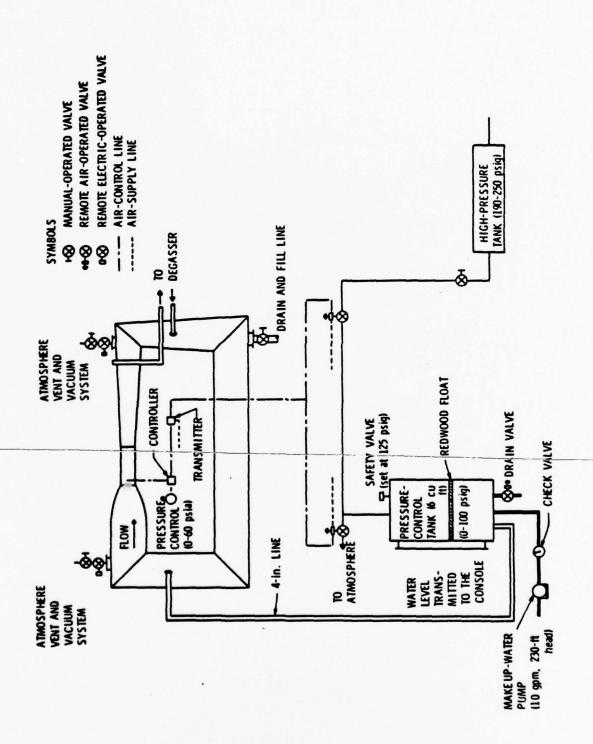


Fig. 3

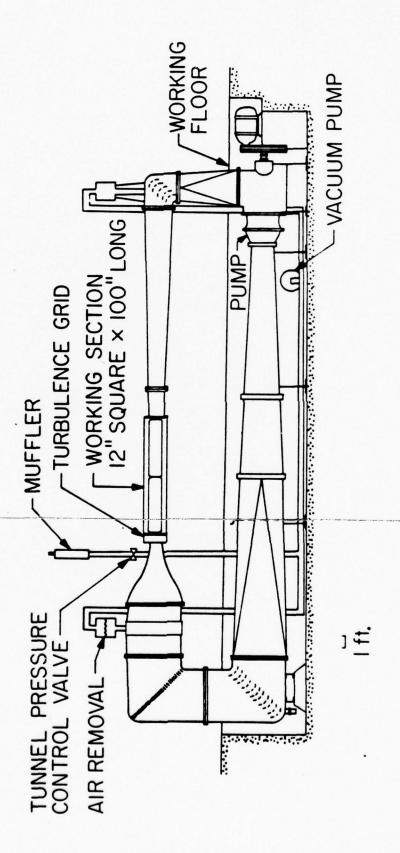


Fig. 4

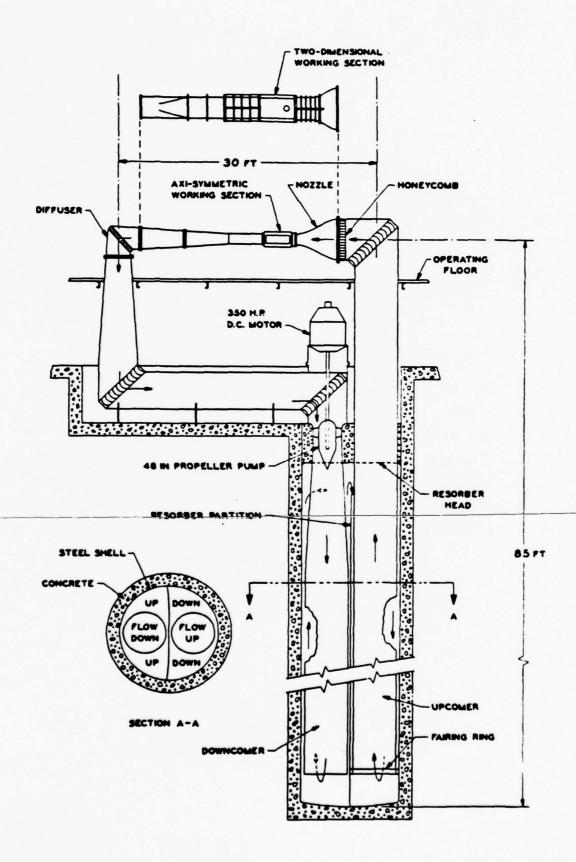
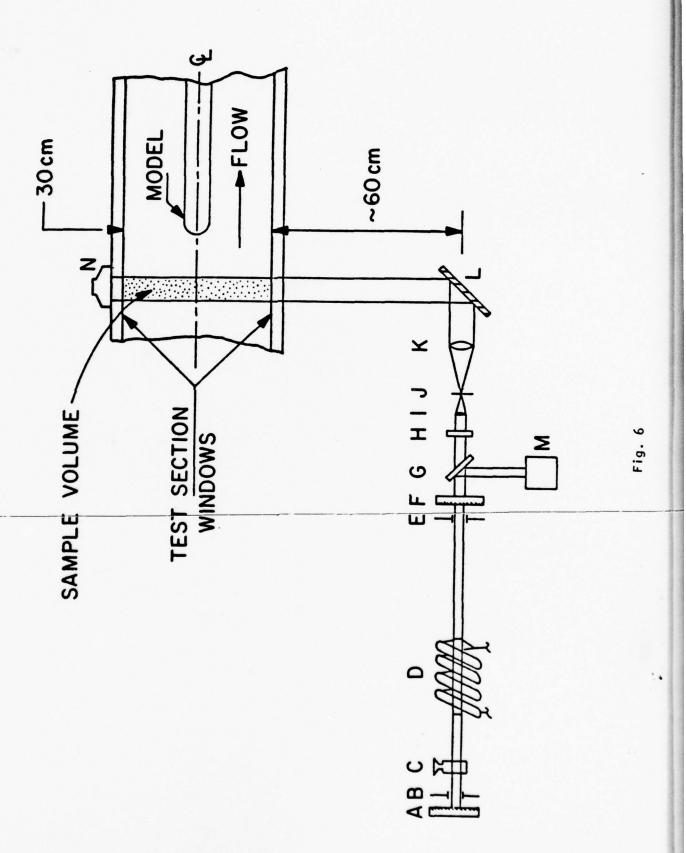


Fig. 5



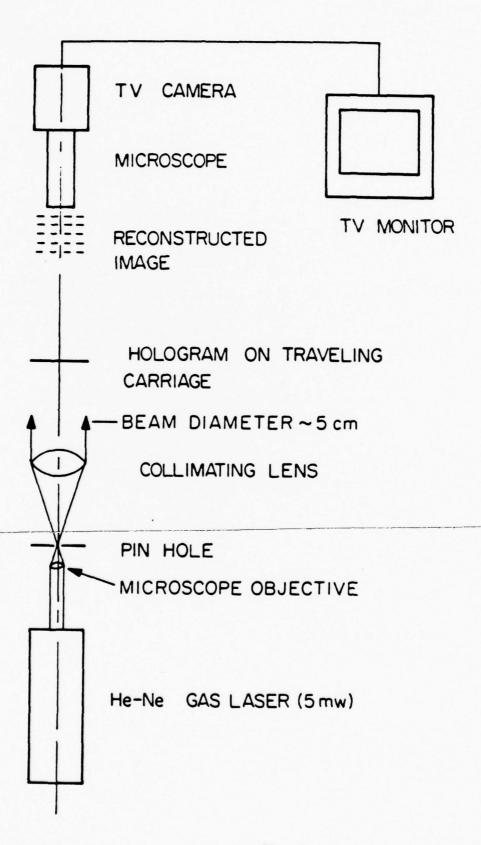


Fig. 7

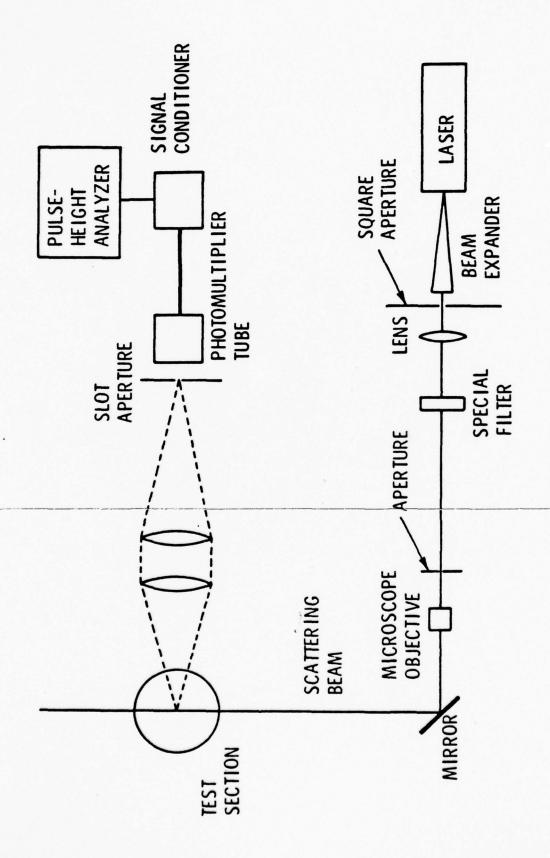
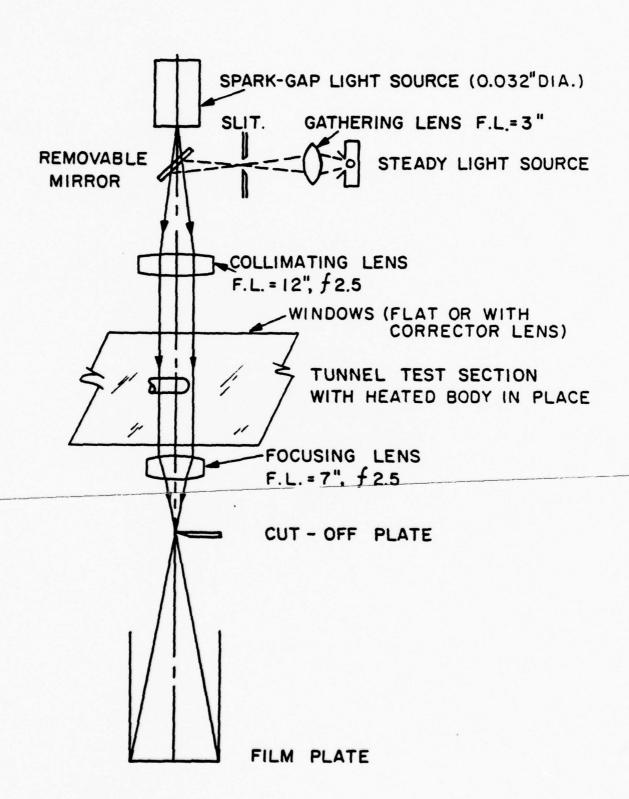
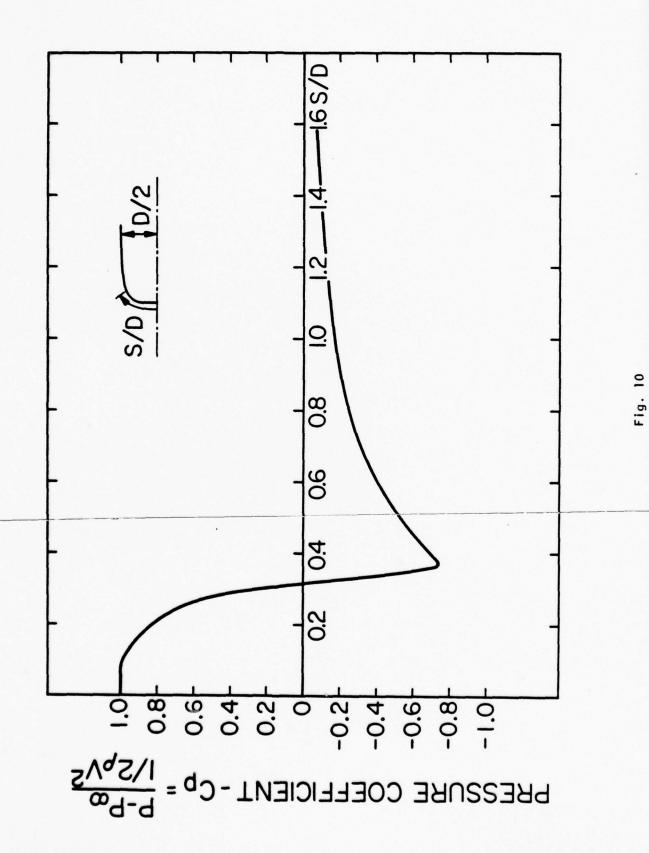


Fig. 8





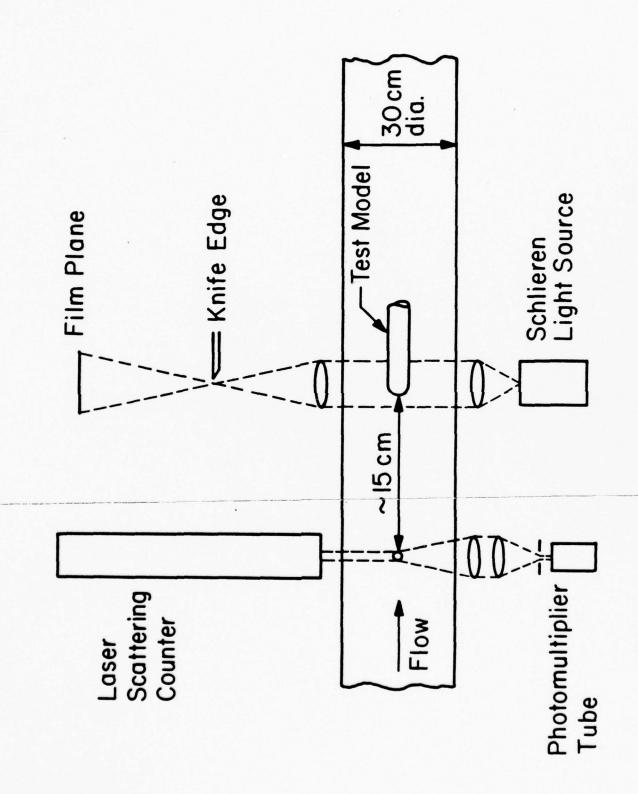
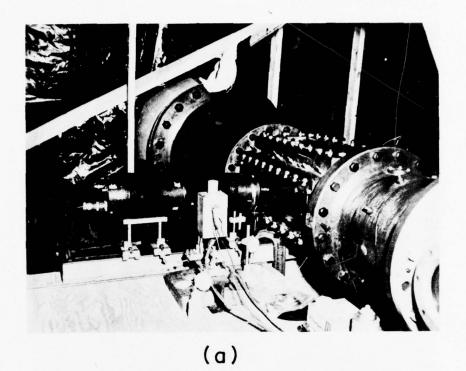
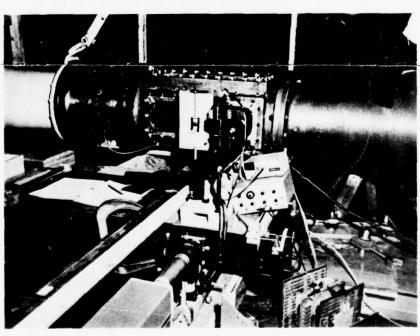


Fig. 11





(b)

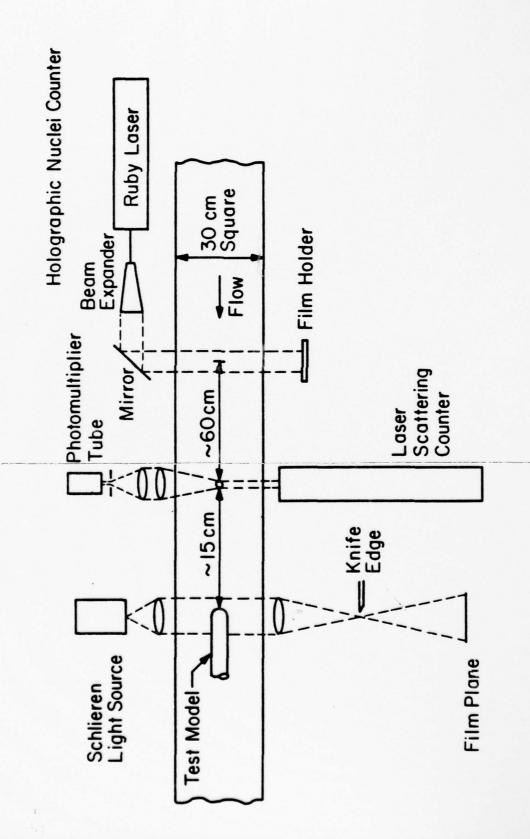
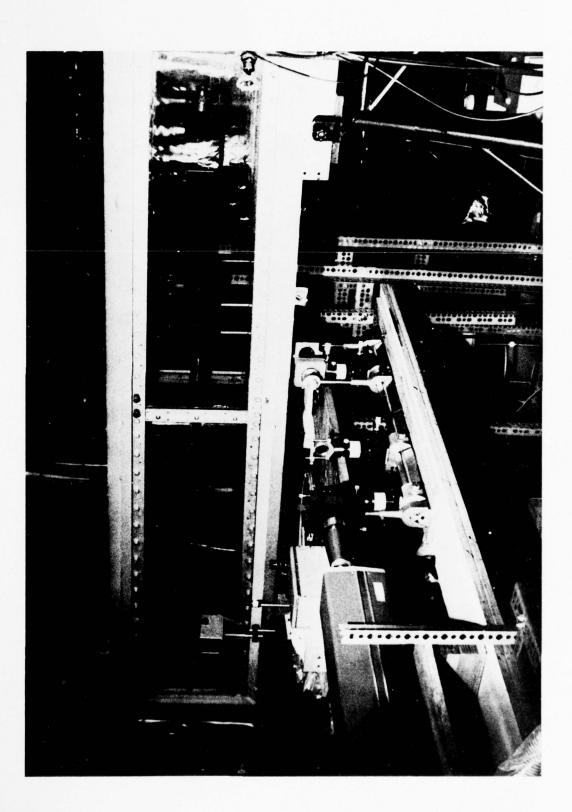


Fig. 13



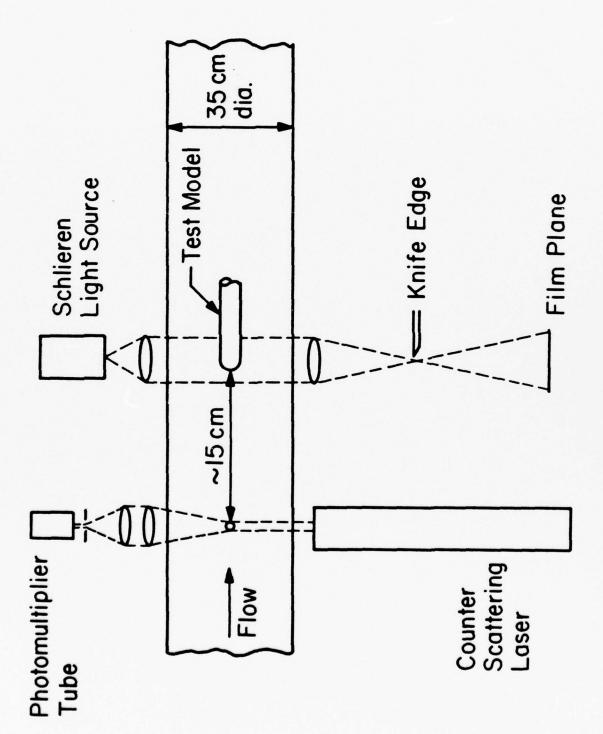
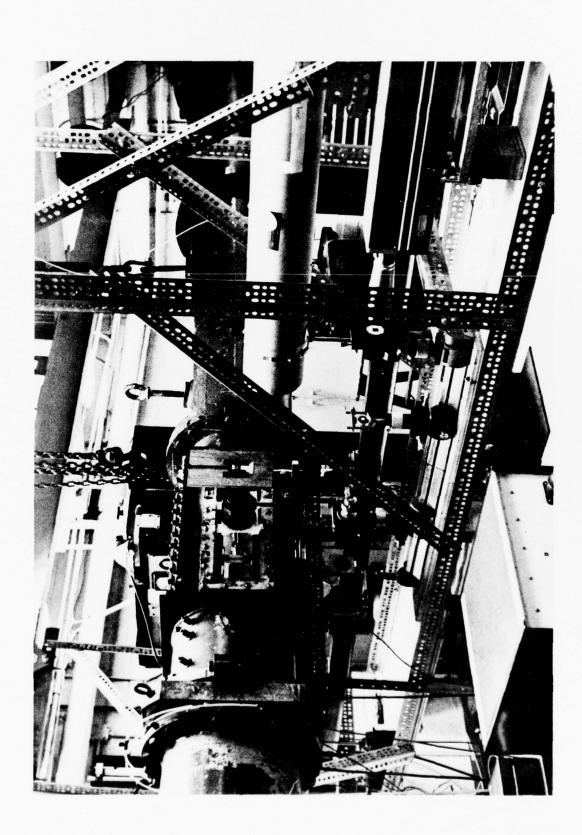


Fig. 15





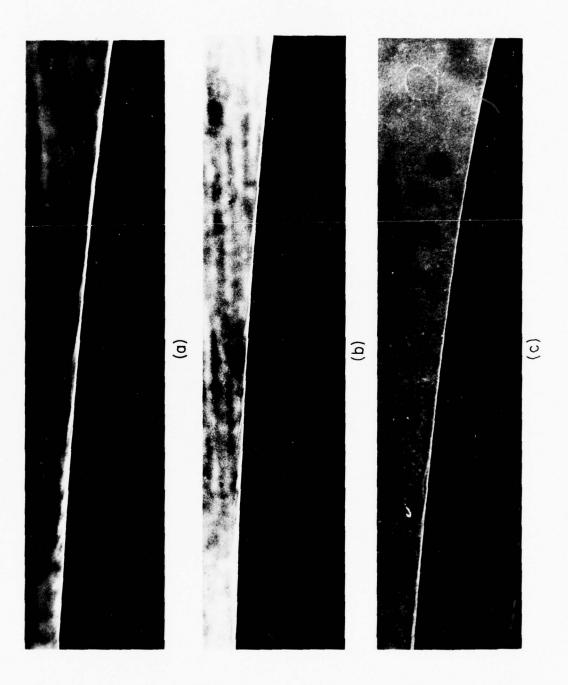
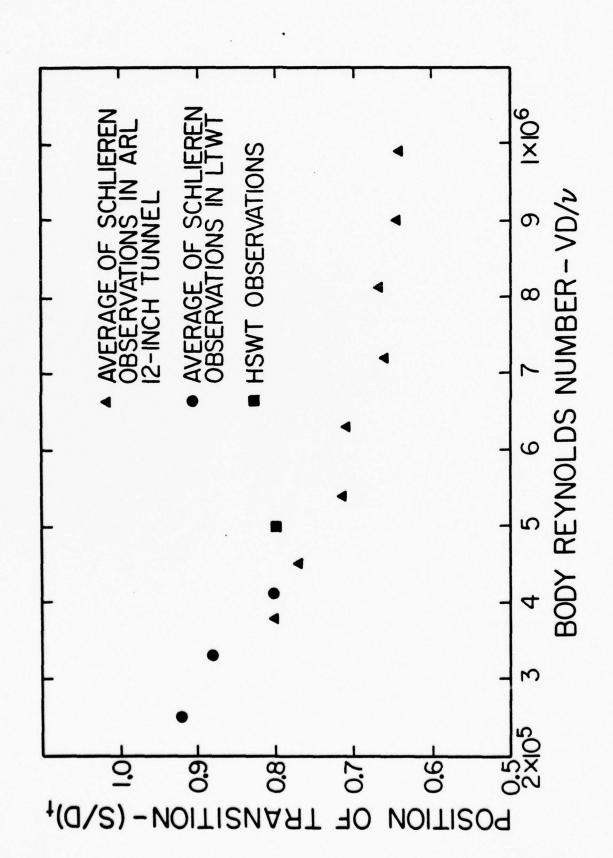


Fig. 17



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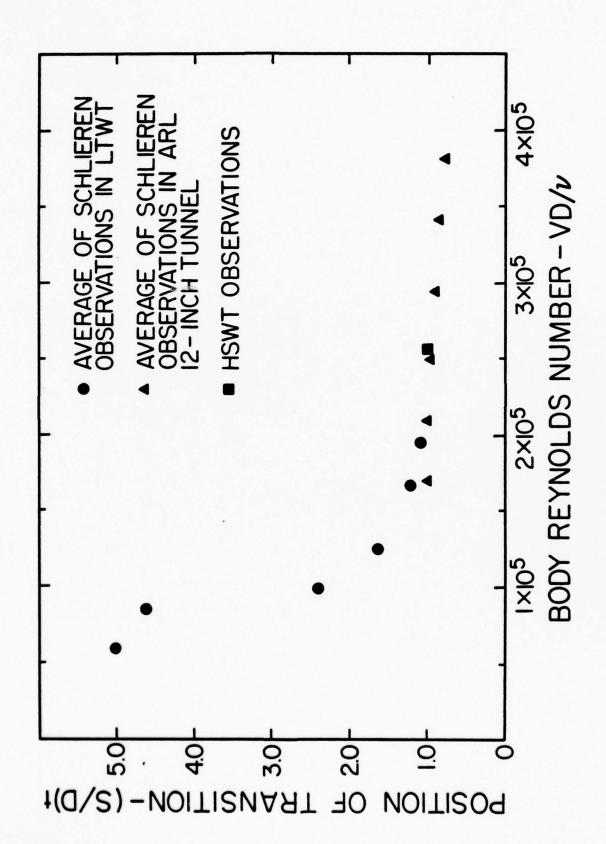


Fig. 19

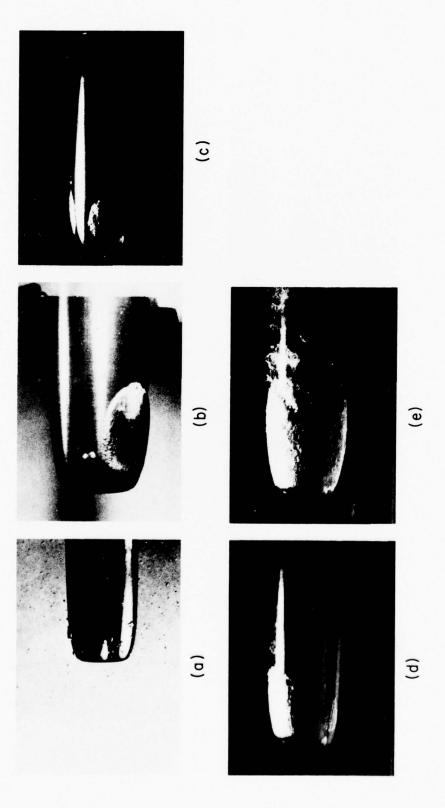
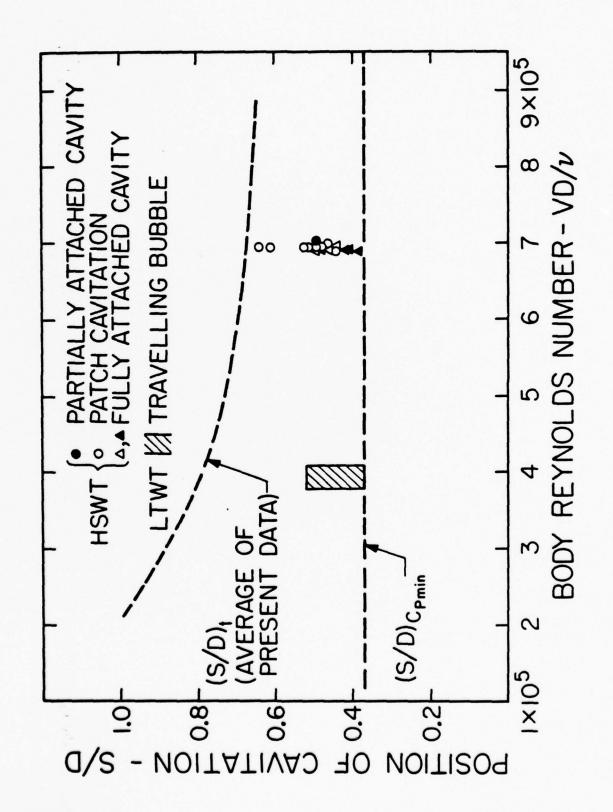


Fig. 20



59

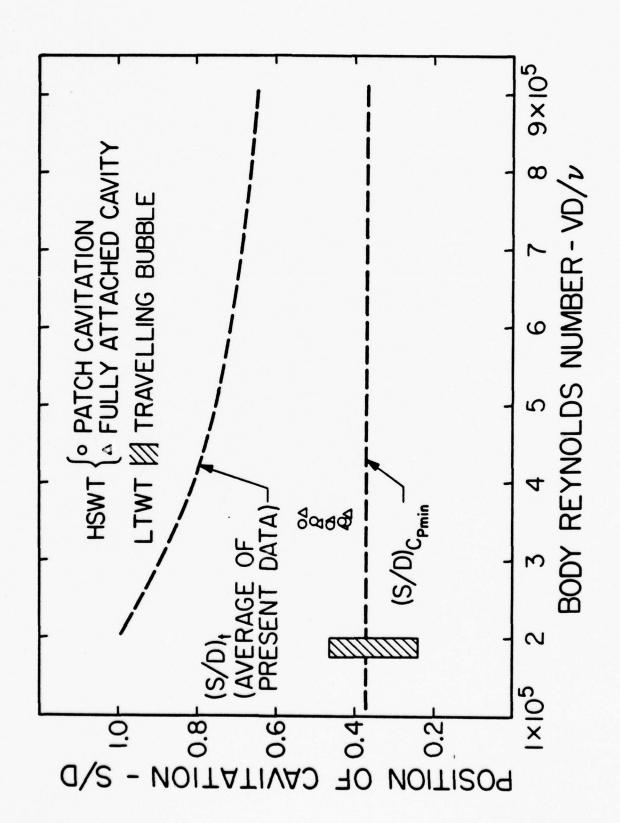


Fig. 22

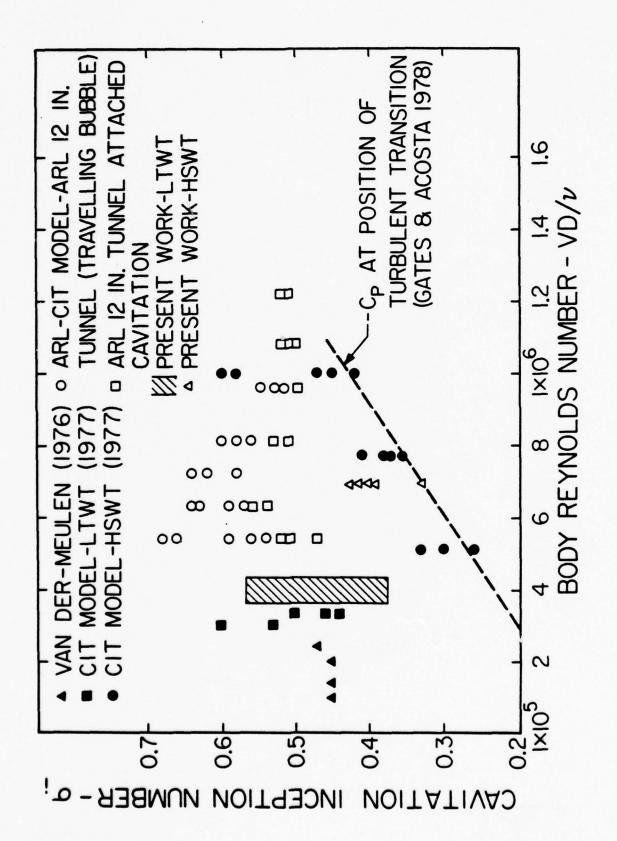


Fig. 23

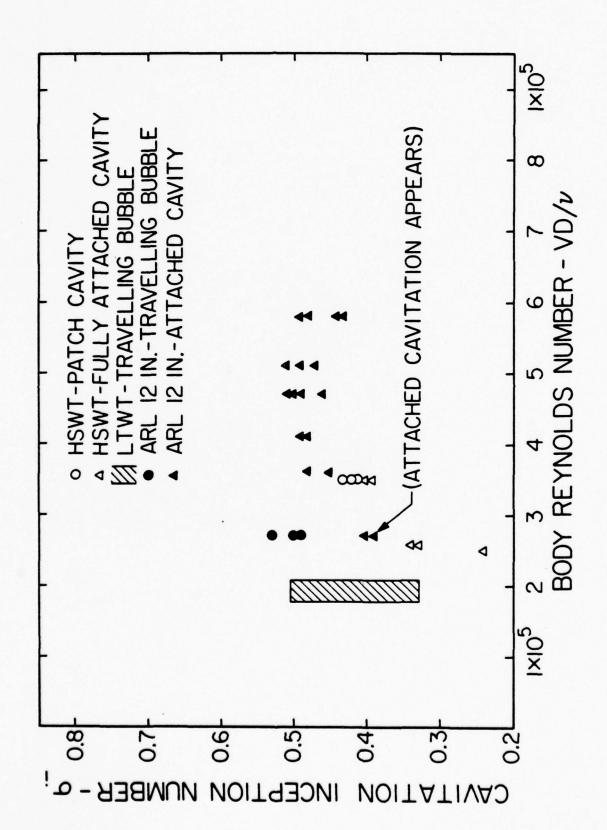
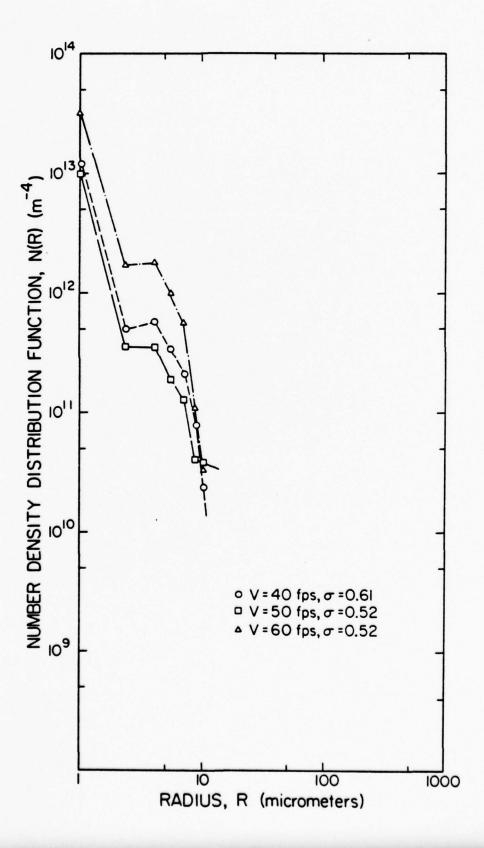
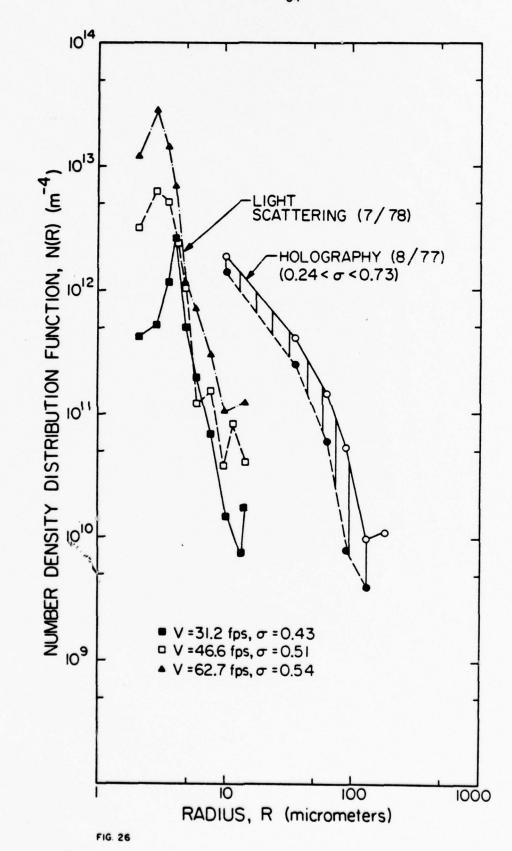
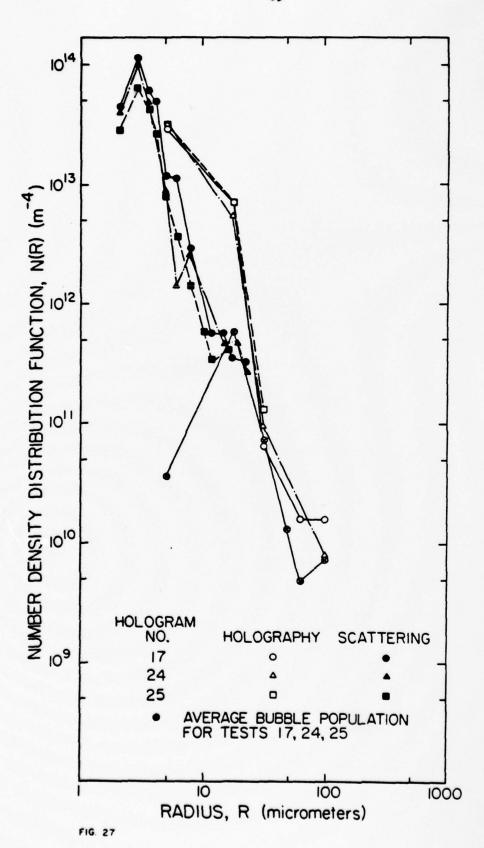


Fig. 24







### APPENDIX 1

#### COORDINATES OF THE SCHIEBE MODELS

The original non-dimensional coordinates for the  $-C_{p_{min}} = 0.75$  Schiebe model were obtained from a report by Schiebe (1972). The number of coordinates was increased by using a spline routine. These coordinates were then substituted into a Douglas Neumann program to determine the pressure distribution. In the appendix we have summarized all this information by including a copy of the potential flow program output. The following symbols used in the program output will be of interest:

- x coordinate measured along the axis of symmetry in inches.
- Y coordinate measured perpendicular to the axis of symmetry in inches (i.e. local body radius)
- Cp- pressure coefficient,
- SUMDS coordinate measured along the surface of the model from the stagnation point in inches.

Ignore all other outputs. The flow program output starts on Page 74.

## APPENDIX II

## CALIBRATION OF THE LIGHT SCATTERING NUCLEI COUNTER

The LSC is calibrated by injecting a solution containing polystyrene spheres of a known size into the sample volume of the counter. The intensity of the scattered light for each sphere diameter is recorded and a calibration of signal amplitude versus particle diameter is obtained. From this calibration then voltages corresponding to the desired ranges are determined and used to program the processor to size and sort detected particles into a series of 16 bins.

The calibration curve of signal amplitude versus particle size for the present LSC system is given in Fig. IIA and tabulated below are the present tests.

TABLE IIA

BIN SIZES FOR TESTS AT ARL\*

Channel	Particle Diameter (micrometers)
1	1 - 3
2	3 - 6.5
3	6.5 - 10.0
4	10 - 12.5
5	12.5 - 16.0
6	16.0 - 19.0
7	19 - 21.5

<sup>\*</sup>Information for channels 8 - 16 is not available. However since more than 95% of all particles were counted in the first 7 channels the loss of this information is not considered serious.

TABLE IIB

Bin Sizes for Tests at Caltech

Particle Diameter (micrometers)
3.5 - 5.0
5.0 - 6.5
6.5 - 8.0
8.0 - 8.5
8.5 - 11
11 - 13
13 - 18
18 - 21.5
21.5 - 24.5
24.5 - 28
28 - 31
31 - 36
36 - 39
39 - 50
750

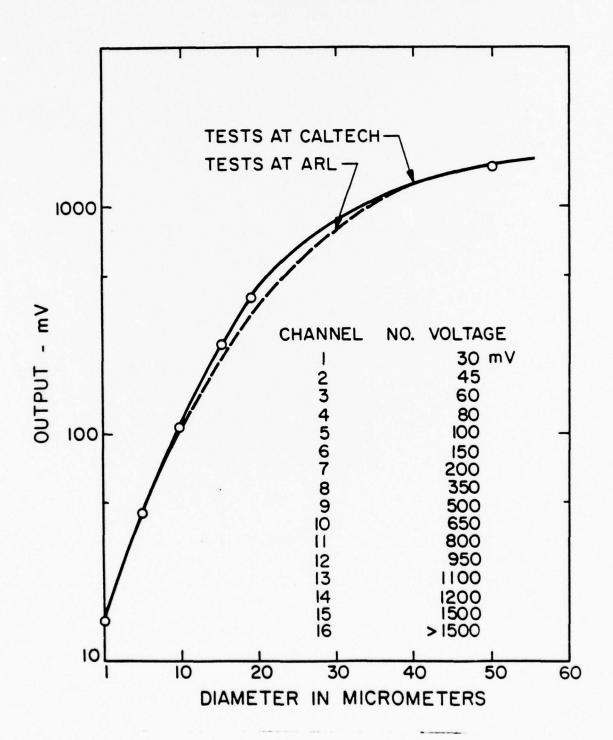


Fig II-A Calibration curve for the LSC showing output amplitude (voltage) versus particle diameter.

### APPENDIX III

#### METHOD AND INSTRUMENTATION

Nuclei counting by holography consists of two processes. In the first one, a sample volume is illuminated by a collimated beam of coherent quasi-monochromatic light causing interference between the main coherent beam, and the light diffracted from the particles in the sample volume. The result is recorded on a high resolution film. In the second process, the photographic record is illuminated by another collimated beam of quasi-monochromatic light, producing a three dimensional image of the original volume. Using a TV vidicon and a monitor, it is possible to size and count the particle distribution.

### Holocamera

The holocamera consists of a light source, a beam expander and a recording film. The coherent light source is a "Q Switched" pulsed ruby laser. A 3 inch long X 0.25 inch diameter ruby is excited by a helical Xenon Flashlamp, which is activated by discharging through it a 1000 Joule pulse in a period of 1.5 milliseconds. The ruby is located in an optical cavity created by 2 flat mirrors. The back mirror is 100% dielectric reflecting surface, while the front (output) one is a single layer Saphir Etalen (reflectivity - 60% in the red light wavelength) from which the output light is emitted.

When the lamp is flashed it excites the ruby rod, which then emits part of the energy absorbed by it. The emitted light oscillates in the cavity, bringing (together with the added energy from the flash lamp) the ruby to an increasingly "unstable state". When the ruby reaches a certain critical state, it emits a giant pulse, part of which is emitted outside of the cavity through the front mirror. Since the duration of one light pulse

is approximately 50 nano-seconds, the same process repeats itself several times during a single flash of the Xenon lamp. Therefore, needing only one single pulse, it is necessary to "Q switch" the laser. The present system is switched by inserting in the cavity a coated surface quartz bottle containing acetone and cryptocyanine. The material in the bottle absorbs part of the light oscillating in the cavity and thus reduces the gain due to these oscillations to a level where only one light pulse is emitted during the operation of the lamp.

The output light passes through a "beam splitter" in which part of the beam is transferred to a pin diode. The diode signal is then displayed on an oscilloscope. The rest of the beam (0.6943 micrometer wavelength) enters a beam expander-spatial filter (10 micrometer pin hole) and after being collimated it passes through the water tunnel to the recording film (Agfa-Gaevert 10E-75) located on the opposite side of the test section.

## Reconstructing System

The reconstructing system uses a Spectra Physics 5 mw Helium Neon Laser as the illumination source. The beam is again expanded, passes through a spatial filter (pin hole) and is collimated by a collimating lens. The collimated beam (2.5 inch diameter) passes through the hologram, and creates a three dimensional image of the original sample volume. By positioning a microscope objective in the field of the reconstructed hologram the image is magnified and focused onto a silicon vidicon. The image is then displayed on a TV monitor. By changing the distance between the objective and the vidicon, one can control the magnification of the nuclei-image shown on the monitor. The hologram itself is mounted on a x-y-z Vernier carriage, making it possible (with the help of scales in all three dimensions) to count the particles in any desired volume.

By careful direct observation of the monitor screen, on which a calibrated reticle is placed, one can count the number and size the particles and bubbles recorded on the hologram. The magnification used is 220%. This magnification makes it possible to observe 10 particles, but the background noise of the system prevents sizing of smaller ones. Higher magnification is impossible due to significant decrease in resolution when the hologram is further magnified.

Since the accuracy of the result depends also on the quality of the hologram and the personal judgment of the one who counts the particles, several holograms were read by different people at different occasions. Since the results were almost equal, it seems that "human error" plays a minor role.

The volume studied is 2.5 cm X 1 cm X 1 cm starting 4 inches from the tunnel window. Since the TV screen covers an area of 1mm<sup>2</sup> at a time, it is necessary to observe 100 different squares, continuously varying the depth of each until the entire volume is covered.

By applying this technique it is possible to distinguish between a bubble and a solid particle. While the particle has got an arbitrary shape, the cross section of a bubble is circular (or at least almost circular). Therefore, distinction is possible only when the image is large enough to observe its shape — in the present magnification, when the size of the bubble/particle is above 20 micrometers.

# Results

The holograms counted were obtained during experiments in the LTWT and HSWT at Caltech. The number of particles and bubbles, found on each hologram is presented in Tables II, V, VI and an approximated number density distribution function:

$$N\left(\frac{R_1+R_2}{2}\right) = \frac{\text{Number of nuclei within radii between } R_1 \text{ and } R_2 \text{ per unit volume}}{R_2-R_1}$$

for each hologram is presented in Figs. 26 and 27.

The results are divided into two groups. The first one refers to the total number of solid particles and bubbles, and the second one refers only to bubbles. As noted before it is impossible to distinguish between bubbles and particles when their size is less than 204, therefore all nuclei found in that region are categorized as particles.

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COMPANY 510N	WITH PARLS		11	0.51608908	0.50743014	0.62035447	0.67996633	0.75819165	0.84923410	00318.000.0	0.92432300	1.03768635	1.17403793		1.25581169	1.29819489	1.30376148	1.33861664	1.30762462	1.30 364132	945:1700-1		1.28948116	1.28164959		7766471701	1.20260669	1.25315762	1.24849606	
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DUUGL AS LUNG	- 1,0125 INCH	LEYMMETAIC FLOW		0.22306001	0.23750301	0.2563.1999	0.26279999	0.28305495	0.29555500	0.30300001	0.311-7999	76997 FIE 0	0.33149999	0.33008003	0.34105003	0.34701003	0.35250002	0.35500002	0.36250001	0.307.0001	0.37640000	0.37530000	0.37750000	0. Jb2+5993	0.38499999	0.3090500	0.39245998	0.39745998	0.39999999	
	Anteus DODY	UN-BUDY UNIFURM AXISYMMETAIC TAANSFURMED COURDINATES	*	0.01272000	0.0147030	0.01795000	0.01900000	0.02156000	C.0266v500	00036750.0	0.0316000	0.02498000	0.04133499	0.04451000	0.04768997	*********	0.00774498	0.00023000	0.00010000	0.07589596	0.07941997	0.03716002	0.09141999	0.10029495	D.V.2994	0.11505002	0-12054948	0.13203496	0.13861998	
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			SIGMA	-0.01586213	-9.01490453	-0.01426687	-0.01342769	-0.01267901	-0.01198050	-0.01119955	-0.01060363	-0.00994387	-0.00943518	-0.00976335	-0.03818425	-0.20762070	7096.0700.9-	-0.00657257	-0.00611241	-0.0360377	-6.60521747	- 9. 07481605	-1.00414045	-0.00342641	-0.00273839		
			CUS A	0.93595	0.94074	37.447.0	61646.0	0.95367	0.45643	C. 36025	6.96.336	0.90641	63696.3	0.97184	1.97433	10976.0	0.97894	0. Энони	0.03274	C . 984 51	C.586.11	C-93759	0.98973	69760.0	87 466.0		
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	CASE NO.		СР	-0.54523087	-6.53144560	-0-51874! 42	-6.50696748	-0.49472237	-0.46227692	-6.46926086	-6.45047163	-6.44402395	-0.43232586	-6.42130375	-C.40h #R514	-0.39063887	-0.35426610	-0.37101445	-0.359111.03	-6.34051500	-0.33346329	-6.32195063	-6.30045601	-0.27537335	-0.24870205	UALT THE THE	0 70 DDG 07 C
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AIRCHAFT CUMP			Іна	-0.63745993	-0.03526678	-0-61106801	-0.63085479	-0.62862813	-0.62038426	-0.62412685	-0.02185585	-0.61956596	-0.01725891	-0.61492461	-0.01250582	-0.61313169	-0.60777050	-6.00533100	-0.00280355	- 0.00036200	-0.59782875	-0.59525943	-0.55127527	-0.53566532	-0.58024460	DINI	
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	SCHIEBE BODY	ON-BODY UNIFORM AXISYMMETMIC THANSFORMED COURDINATES	*	0.14437997	0.1510250	0.15736060	0.10516003	0.17272002	0-10470496	0.19312000	0.19742300	0.21109998	0.21573999	0002200700	0.24132938	0.24675500	0.0010000	0.27637594	0.24267994	0.20296001	0.31742001	0.32270301	0.33784030		0.41461303		. 365
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DUUGLAS AIRCRAFT COMPANY LONG SEACH DIVISION

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5-10600038	0.524.99999	68650899*0-	0.64610057	C.28411388	0.24254	0.97014	-0.02331039	-0.00000226
2 1 4 4 9 5 3 6 2	0.53750002	-0.467646748	0.87277615	C - 234261 42	0.24253	6.97014	-0.02262508	-0.000000-0-
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######NING### UUTBUL FIELD WIDTH 13G SMALL. CONDITION OCCURRED DUFING A FORMATTED #HITE ON FORTEAN UNIT 6 WHICH IS ATTACHED ID \*\*SINK\*, THE WRITE IS SEQUENTIAL AT RECORD NUMBER 733, FOR THIS AND ALL FUTURE OCCUPAENCES OF THIS CONDITION, A

FIELD OF \*\* S WILL BE WRITTEN.

ADDED MASS = 2.15955257 VOLUME = ##########

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ווף רוצ		×>	0.936506.99	6.57435051	0.58205251	0.93466212	C. SH575449	0.99626190	ECUTION.	*************			AL DISPLACEMENT
, ,	۳ د د ه	IHd	-0.51748586	-6.47872937	-0.4576 3900	-0.44125748	-0.426537HZ	-0.4125766n	END OF FILE ON #SJUKLE#CAUSES TERMINATION OF EXECUTION. FORMAI (10A6,2%,A9,8%,[4/27]1,4%%,14)	****************	Dus.	.45x.141	SUBPHUGNAM THACEBACK. 10 EKNUR JCCUMKED IN HOUTINE PARTI AT HEXADECIMAL DISPLACEMENT +57C CALLEU FRUM HUUTINE MAINELANIS CALLED FRUM "SYSTEM"
	AXI SYMMETRIC	>	0.0	0.0	0.0	0.0	0.0	0.0	J.C.E & CAUSES T	***********	L UUTFUT FOLL	AU. 3X, 14/2711	IN ROUTINE PA
HILD BUDY - 1 125	OFF-BUDY UNIT-JAM AXISYMMETRIC FLUM THANSFURMED COUNDINAIES	×	-1-000000000	-4.00000000	-3.30000000	-4.00000000	= 300.0000 C. C.	-6.36600000	END OF FILE ON #SJURCE#CAUSES TERMINATI FORMAI (1646,2%,40,8%,14/2711,45%,14)	***********	BATCH FURFRAN DEBUG GUTPUT FOLLDES.	FUKMAT (10A0.2X, 40, 4X, 14/2711, 49X, 14)	SUBPRUGNAM TRACEBACK.  10 ERRUR JCCURRED IN ROUTINE P. CALLEU FRUM ROUTINE MAINEISNIS CALLED FRUM "SYSTEM"
	0FF-16		-	. ~	1 17)	•	4	3	FUND OF	*****	ВАТСН	FURMAT	SUBPRU 10 EKK CALLCU CALLCU

DOUGLAS AIMCHAFT COMPANY LONG DEACH DIVISION

UNIT 5 15 ASSIGNED FOR #SGUNCE# UNIT 6 15 ASSIGNED TO #SINK# UNIT 10 15 ASSIGNED TO -TIC UNIT 11 15 ASSIGNED TO -TI1 UNIT 12 15 ASSIGNED TO -TI2 UNIT 15 15 ASSIGNED TO -TI2 UNIT 16 15 ASSIGNED 10 -115
UNIT 16 15 ASSIGNED 10 -116
GUSER 15 ASSIGNED TO #MSCUNCE+
SCARDS 15 ASSIGNED TO #MSCUNCE+ SEKCUM IS ASSLUNCU TO BHOLSE SPRINT 15 ASSLUNED TO \*MSINK \* UNIT 14 IS ASSIGNED IL -114 UNIT USAGE TALLE

ERRUR NUMBER=130

SEVERITY LEVEL=3

FORTIAN TO EMADE CAUSES A METURN TO THE SYSTEM. OCISOLIS 14.0CB RC=10

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												D ALPHA		0.0	0.0	0.0	0.0	-0.30774099	-0.49439533	-0.48111844	-0.38932073	-0.31325972	-0,20069826	-0.26424628	-0.36593419
			section.									SUMDS	0.00728000	0.01455000	0.02183000	0.02911000	0.03638000	0.05500026	61500219	0.10000843	70610521-0	0-15041423	0-17505211	0-30007697	0.22510517
		ON BOAD	Test Sec									DELTA S	0.00728000	0.00727000	0-03728000	0.00723000	0-0727000	0.01862026	0.02000194	0.02500624	0.02501068	0-02534520	0.02463793	0 02502257	0.02503021
			" 0/A				PANY		HY = 0.0			COS A	0.0	0.0	0-0	0.0	0-0	0.00537049	0.01399863	0.02239441	0.02918752	0.03465221		0.047740.0	0.04914067
AND CROSSFLOW	:	S II WALLS	12		· man make the line	SURFACE OF REVOLUTION OFF-BODY POINTS MAT CORRESPONDENCES FOR SOLUTION	AIRCRAFT COMPANY BEACH DIVISION	WITH WALLS	0.0	0.0	DRMEDI	SIN A	1.00000000	1.000000000	1.00000000	1.00000000	1.00000000	0.9998599	78200866.0	0.9976942	0.9957412	0.99939954	0-98927211	0.99908543	0.99879205
AXISYMMETRIC AND CROSSFLOW	***** CASE CONTROL DATA ****	S INCH - UTTH USILS				SURFACE OF REVOLUTION OFF-BODY POINTS	DOUGLAS	2.025 INCH - W	1	YE	COORDINATES (UNTRANSFORMED)	>	0.00	0.00728000	0.01455000	0.02183000	0.02911900	0.03638303	0.05500000	0.07499999	0.09999996	0.12500000	0.15037996	0.17499995	0.19599999
PROGRAM 500	***** CAS	SCHIFBE BOOK	8001ES = 2	MACH NO.= 0.0	A CONTRACTOR OF THE PARTY OF TH	SURFACE OFF-BODY		SCHIEBE BODY		XF . 0.0	ON-BODY COORDIN	×	0.0	0.0	0.0	0.0	0.0	0.0	0.0010000	0.03038000	0.00394000	0.00167000	0.03255000	0.00349000	0.00456000
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	0.25703824	0.27519596	0.29483378	0.32742441	0.35043234		0.37554729	0.40068835	0.42585826	0.44716263	0.47632366	0.50144143	7663746	FEEE 0176-0	0.55255431	0.58569092	0.60416871	0.62838227	0.66030586	0 485 20002	26625363	27270712	18181811-	0.73344457	0.74863911	0.76482713	0-78210503	76035000	4.80033334	0 04133133	0.84132123	0.86388119	0.88806295	0.91399825	0.92768425
	0.03193307	0.01815772	0.01963786	0.03259065	0.02300798		0.02511497	0.02514109	0.02516992	0.02130437	0.02916107	0.02531601	0.0350	0-02801112	10106520-0	0.03313663	0.01847781	9.02421359	0-03192361	0 03600300	D1616150-0	D1806020-0	76688371101	0.01427712	0.01519455	0.01618804	0.01727791	16.737000	0.010705693	- constant	0.02105613	0.02256661	0.02417521	0.02593534	0.01368605
	0.05762048	0.05443548	0.06976306	0.07855016			0.09556049	0.10580295	0.11601138	0-12767339		0.15799029	•			0.23357826	0.27221835	0.31758994		0 40853949	200000000000000000000000000000000000000	101700000	anger na gen	0.71373010	0.75290424	0.78638297	0.8154887								0.93087602
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81	3.16443920	1.00000000					-0.49942166	-
	3.38221931	1.00000000	0.0	1.00000000	0.43556023	4.25460625		
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PENNSYLVANIA STATE UNIV UNIVERSITY PARK APPLIED RESE-ETC F/G 20/4 CAVITATION INCEPTION AND NUCLEI DISTRIBUTIONS JOINT ARL/CIT EXP-ETC(U) SEP 79 E M GATES , M L BILLET , J KATZ N00014-79-C-6043 AD-A077 633 UNCLASSIFIED NL END 2 OF 2 DATE ADA 077633 1 - 80DDC

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						z	-0.00004274	-6.00004464	-0.00004178	-0.00004274	-C.00004274	-0.00004303	-0.00004232	-0.00003821	-0.00003541	-0.00003654	-0.00003737	-0.00003606	-0.00003308	-0.00003290	-0.00003356	-C.00003338	-0.00003129	-0.00002927	-0.00002897
						SIGMA	-0.14428157	-0.14422148	-0.14409387	-0.14386898	-0.14343065	-C.14306045	-0.14262742	-0.14266235	-0.14278024	-0.14293545	-0.14295352	-0.14272124	-0.14234239	-0.14208210	-0.14152145	-0.14122152	-0.14098901	-6.14025772	-0 13071728
						COS A	0.0	0.0	0.0	0.0	0.0	0.00537	0.01400	0.02239	0.02919	0.03%65	0.03815	0.04276	0.04914	0.05762	0.06444	0.06976	0.07855	0.08736	799000
				3.		SIN A	1.00000	1.00000	1.00000	1.00000	1.00000	66666 0	0.99990	0.99975	0.99957	0.99940	0.99927	60666.0	0.59879	0.59834	0.99792	0.99756	16966.0	91966.0	6 ,500 0
				CASE NO.		CP	0.99999022	0.99991089	0.99975204	0.99951851	0.99925297	0.99843639	0-99650234	0.99345046	0.98912948	0.98367816	0.97736615	0.96960443	0.96012712	0.94713080	0.93362370	0.92317349	0.90543890	0.88324672	9/30000
ADDY = 0.0			ANY			11	0.00312923	0.00944209	0.01574836	0.02194327	0.02733194	0.03954294	0.05914108	0.08092928	0.10426188	0.12775695	0.15044558	0.17434335	0.19568200	0.22993308	0.25724757	0.27717596	0.30750793	0.34169185	20030000
0.0 0.0 0.0 EDRMED1			AIRCRAFT COMPANY BEACH DIVISION	WITH WALLS		PHI	-1.07514954	-1.07510281	-1.07501030	-1.07487166	-1.07468700	-1.07428837	-1.07351112	-1.07235050	-1.07069016	-1.06857872	-1.06600952	-1.06297267	-1.05942535	-1.05482769	-1.05021954	-1.04643250	-1.04074287	-1.03398991	. 0175016.
3 MX = 0.0 0.0 ADDX = 0.0 0.0 YE = 0.0 CORDINATES_(UNIRANSEORMED)	Y-0FF	0.00	DOUGLAS	2.025 INCH - 1	ISYMMETRIC FLOW		0.00364000	0.01601500	0.01819000	0.02547000	0.03274497	0.04568997	0.0649994	0.08749998	0.11249995	0.13768995	0.16268992	0.18749994	0.21249998	0.24093997	0.26593995	0.28479499	0.31093494	0.33853996	0.3499956
NN = 3 THETA = 0.0 XE = 0.0 DEF-BODY CORROLA	x-0FF	-2.00000000	,	SCHIEBG BODY	ON-BODY UNIFORM AXISYMMETRIC TRANSFORMED COORDINATES	×	0.0	000	0.0	000	000	0.00005000	0.00024000	0.00066000	0.00130500	0.00211000	0.00302000	0.00402500	0.00517500	0.00671000	0.00821500	0.00948500	0.01145000	0.01373500	0004141000
9		- 2 -			0N-8		1	, ,							2	:	:	:	:		:	:	:		61

	DY UNIFORM AXISYMHETRIC FLOW TRANSFORMED COORDINATES  0.02544000 0.44612998 0.02544000 0.44612998 0.02544000 0.44612998 0.0355000 0.47499996 0.03585500 0.51279974 0.03136500 0.5255996 0.04313000 0.5255996 0.04313000 0.5255996 0.04313000 0.52559996 0.05398999 0.56110975 0.05590000 0.5999996 0.0538899 0.61147976 0.06358999 0.61147976	0.43556494 -1.00448227 0.  DOUGLAS AIRCRAFT COMPANY LONG BEACH DIVISION LONG BEACH DIVISION  2.025 INCH - MITH MAILS  0.44612998 0.44612998 0.46056497 -0.99523491 0. 0.46056497 -0.99523491 0. 0.51279996 0.51279974 -0.97319907 0. 0.51279974 -0.97319907 0. 0.5255996 0.52799996 0.59110975 -0.93169200 0. 0.59999996 0.5999996 0.59110975 -0.91873473 0. 0.63759999	0.44259203 0.47954732 PANY 0.52096581 0.62665033 0.68668121 0.85813057 0.93325377	CASE ND  CAS	0.99325 0.99182 0.99002 0.98417 0.97234 0.96224 0.96224 0.96224	0.11601 0.12767 0.12767 0.15799 0.17723 0.19959 0.27722 0.39759	-0.13810199 -0.13693804 -0.13597041 -0.13513193 -0.13213193 -0.13248674 -0.12646847 -0.11588455	-0.00002903 -0.00002474 -0.00002253 -0.00002331 -0.00002331 -0.00002331 -0.00002331 -0.00002283
	0.65223998 0.66299963 0.67376000 0.68209457	-0.86497611		-0.40659046	0.86104	0.50854	-0.08704203	-0.00002313
0.11038995 0.101038995 0.011548495 0.012057996 0.012057996 0.013202000 0.012202000 0.0122000 0.0122000 0.01220000 0.0122000 0.0122000 0.0122000 0.0122000 0.0122000 0.01220000 0.0122000 0.0122000 0.0122000 0.0122000 0.0122000 0.0122000 0.0122000 0.0122000 0.0122000 0.0122000 0.0122000 0.0122000 0.0122000 0.0122000 0.0122000 0.0122000 0.0122000 0.0122000 0.0122000 0	69521475 6959999 70499992 7099998 71499568	-0.84506840		-0.74721813	0.70042	0.71373	-0.05957862 -0.05291888 -0.0466G688	-0.000010198
		-0.818753360 -0.81875318 -0.80997592	1.31128554	-0.74535085 -0.73524094 -0.71825796	0.57877	0.84046	-0.03651353 -0.03651353 -0.03244713	-0.00000966
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DOUGLAS AIRCRAFT CCHPANY LONG BEACH DIVISION

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-97 -0.00000626 -0.00000715 -0.00000656 -0.00000572 -C.00000697 -0.0000000--0.00000644 -0.00000554 -C.00000536 -0.00000536 -0.00000507 -0.00000530 -0.00000489 -6.00000447 -0.00000423 -0.00000405 -0.00000292 -0.00000292 -0.00000477 -0.00000507 -C.00000381 -0.0000050 -0.01211775 -0.01142377 -0.00959615 -0.00898720 -C.00837996 -0.01447648 -0.01371105 -0.01020734 -0.00781133 -0.00727487 -0.00496143 -0.00435580 -0.01618297 -0.01527132 -0.01288001 -0.01079564 -0.00676685 -0.00627554 -C.00580623 -0.00538427 -0.00357304 -0.00288084 SIGHA 0.93587 0.96025 0.96915 99056-0 0.94503 C.55683 0.96634 0.97181 0.97433 0.97884 0,98969 0.99208 0.94912 0.95312 0.96340 19976.0 98086-0 0.98275 0.98450 11986.0 0.98760 0.99408 COS 0.35236 0.29065 0.27914 0.24648 0.32700 0.19469 0.33935 0.25728 0.23576 0.21475 0.12560 0.10869 0.30258 0.26806 0.22512 0.18494 0.16612 0.15696 0.14323 16416-0 0.20461 0.17537 SIN -0.57814789 -0.46533489 -0.54011440 -0.52794838 -0.51506710 -0.50206089 -0.48930550 -0.47714520 -0.41616249 -0.30644989 -0.55221462 -0.45343399 -0.44123077 -0.42873955 -0.40357685 -0.39096928 -0.37824154 -0.35351658 -0.2799496 -0.56513977 -0.36565208 -0.33191681 3 1.23410210 1.25105572 1.24101353 1.23088074 1.22558594 1.21537876 1.21051025 1.20051289 1.19529915 1.19002628 1.18472672 1.16861153 1.14300060 1.13136864 1.25624371 1.24587917 1.22037125 1.20558453 1.17939377 1.17398548 1.16340733 1.15408707 F AIRCRAFT COMPANY BEACH DIVISION -0.75439119 -0.63188225 -0.74057388 -0.73589790 -0.73118407 -0.72164315 -0.71681058 -0.71192479 -0.70196855 -0.69174808 -0.68653131 -0.68123585 -0.66482456 -0.64431471 -0.74520838 -0.72643256 -0.70697927 -0.69689274 -0.67585677 -0.67039049 -0.65616387 -0.74981201 2.025 INCH - WITH WALLS H IDY UNIFORM AXISYMMETRIC FLOW TRANSFORMED COORDINATES 0.84249973 0.80749989 0.83249598 0.84745985 0.85249996 0.85749960 LONG 0.81749964 D.8549996 0.86249971 0.86749983 0.87249994 0.87500000 0.87749958 0.88249969 0.89749956 0.91499998 \$666606 · U 0.81249952 0.81500000 0.8199999 0.82249975 D.8249999 0.82749987 0.8299998 0.8349998 0.83749962 D. 8399997 0.8449997 96666648-0 0.8559995 0.86499955 266666980 0.00000000 0.88499939 0.88745981 0.8899999 0.99249992 0.89499998 856666680 0.90499973 15666606.0 0.91999996 SCHIEBE BODY 0.30897999 0.32313496 0.33036000 0.33789498 0.36117995 0.36940992 0,38623995 40382493 0.42219996 0.43158996 0.45124996 46155494 0.47185999 0.48267996 0.49349999 0.55275488 0.56535000 0.57863474 0.60595465 0.66539955 0.71567965 0.75022995 0.82922000 0.30204999 0.31590998 0.34542996 0.35330492 0.37763995 0.39483994 0.41280997 0.50486994 0.54015994 95616165.0 0.63482952 0.64966995 0.68112999 0.78972483 0.52819967 0.61998999 0.51623794 ON-800Y 45 53 9 20 54 59 63 47 3 25 62

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ON-BOOY UNIFORM AXISYMETRIC FLOW	TRANSFORMED COORDINATES
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	Ŧ	-0.24989933	-0.26446462	-0.27882051	-0.29308283	-0.30729926	-0.32149112	-0.33566779	-0.34983146	-0.36397898	-0.37810022	-0.39217371	-0.40615731	-0.41996932	-0.43345767	0 4443469	,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,	-0.4583103	-0.47561347	-0-47940964	-0.47906506	-0.4745944	71898777	AIRCRAFT COMPANY BEACH DIVISION	MITH WALLS		PHI
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D 100C 170.0	0.02407867	0.02558595	0.02646005	0.02695847	0.02724177	0.02740449	0.02750027	0.02756000	0.02760196	0.02763861	0.02768224	0.02775043	0.02787584	0.02813512	0.02874845	0.03083843			CASE NO		٨٨	0.0	0.0		
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